

Crystal field in multielectron 3d metallic ions

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Crystal field

- I. Introduction / Definition
- II. One electron (3d electron)
 - 1. Symmetries
 - 2. Hamiltonian
 - 3. Crystal field splitting of 3d electron
- III. Multi-electrons ion (3d electrons)
 - 1. Hamiltonian
 - 2. Configurations/ spectroscopic terms
 - 3. Symmetry
 - 4. Examples
- IV. *Crystal field and magnetism*
 - 1. Hamiltonian
 - 2. Magnetic moment (spin and orbit)
 - 3. Zeeman effect / paramagnetism

Crystal field theory : origin



Hans Bethe (1906-2005) [2]

(Nobel Prize in Physics 1967)



In 1929 (University of Tübingen) [1]

“ Either we treat the crystal as a complete whole (spatial-periodic potential field and wave function) “ (*F. Bloch theory*)

or

“ one can start from **a free atom and treat its disturbance in the crystal** “



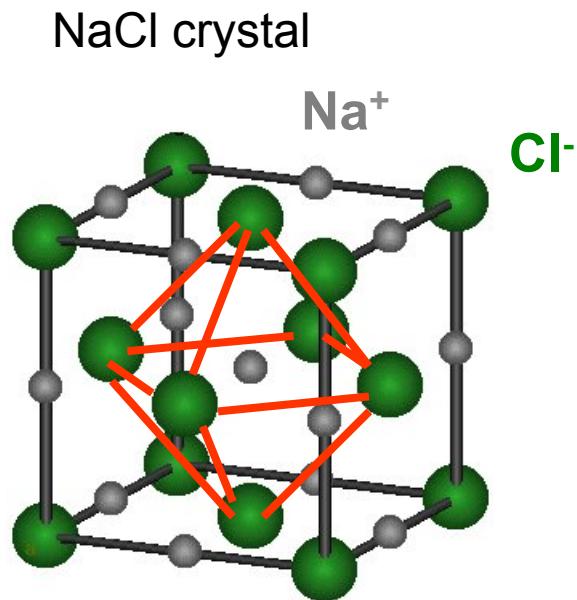
“the atom in the crystal is influenced by the other atoms by an **electric field** of certain symmetry“ = **Crystal field**

[1] H. Bethe, Annalen der Physik, 1929

[2] W. Kutzelnigg, Angew. Chem. 44, 25 (2005)

Crystal field : origin

Hans Bethe (1929) : Model used for NaCl crystal.



Crystal field (CF)

Electric potential generated by the charges of the neighbor atoms

Na⁺ ion surrounded by 6 Cl⁻ ions (nearest neighbors)

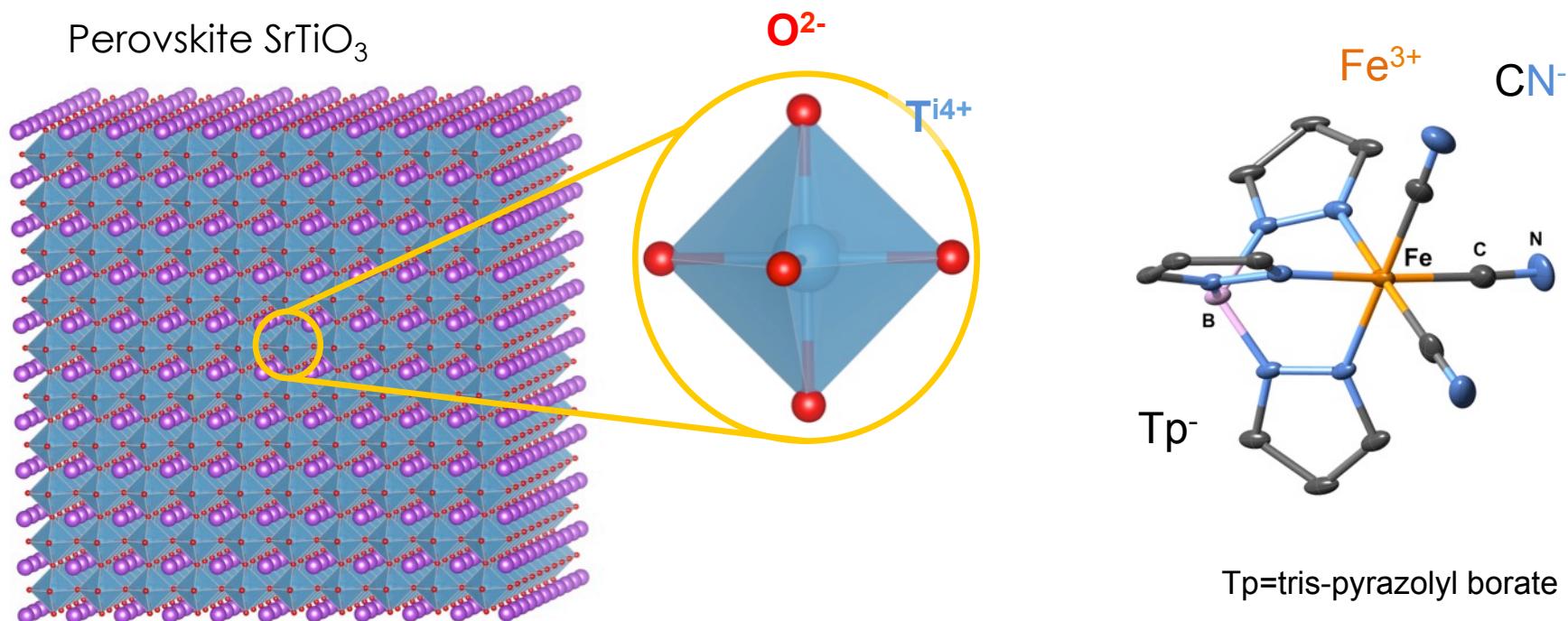
- CF = Electrostatic potential produced by 6 negative charges:

$$V_{CF} = \sum_{i=1}^6 \frac{e}{r_i} \quad (r_i = \text{Na-Cl distance})$$

- Depends on the local symmetry

Crystal field : origin

Analogy in solid-state or coordination complexes :
Consider the ligands as charged spheres



Crystal field model

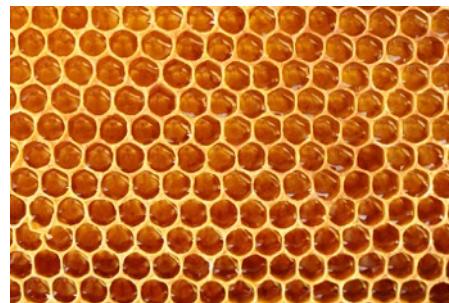
- Local model (restricted to the first neighbors /ligands)
- A purely *ionic* model for transition metal complexes.
- Ligands are considered as point charge.
- Predicts the pattern of splitting of d-orbitals.
- Used to rationalize spectroscopic and magnetic properties.

II.1. Symmetries

Symmetry : responsible for many physical and spectroscopic properties of compounds

Group theory : powerful tool

- simplify calculations,
- predict some properties
- defines the language of labeling



Symmetry operations : some notations to know

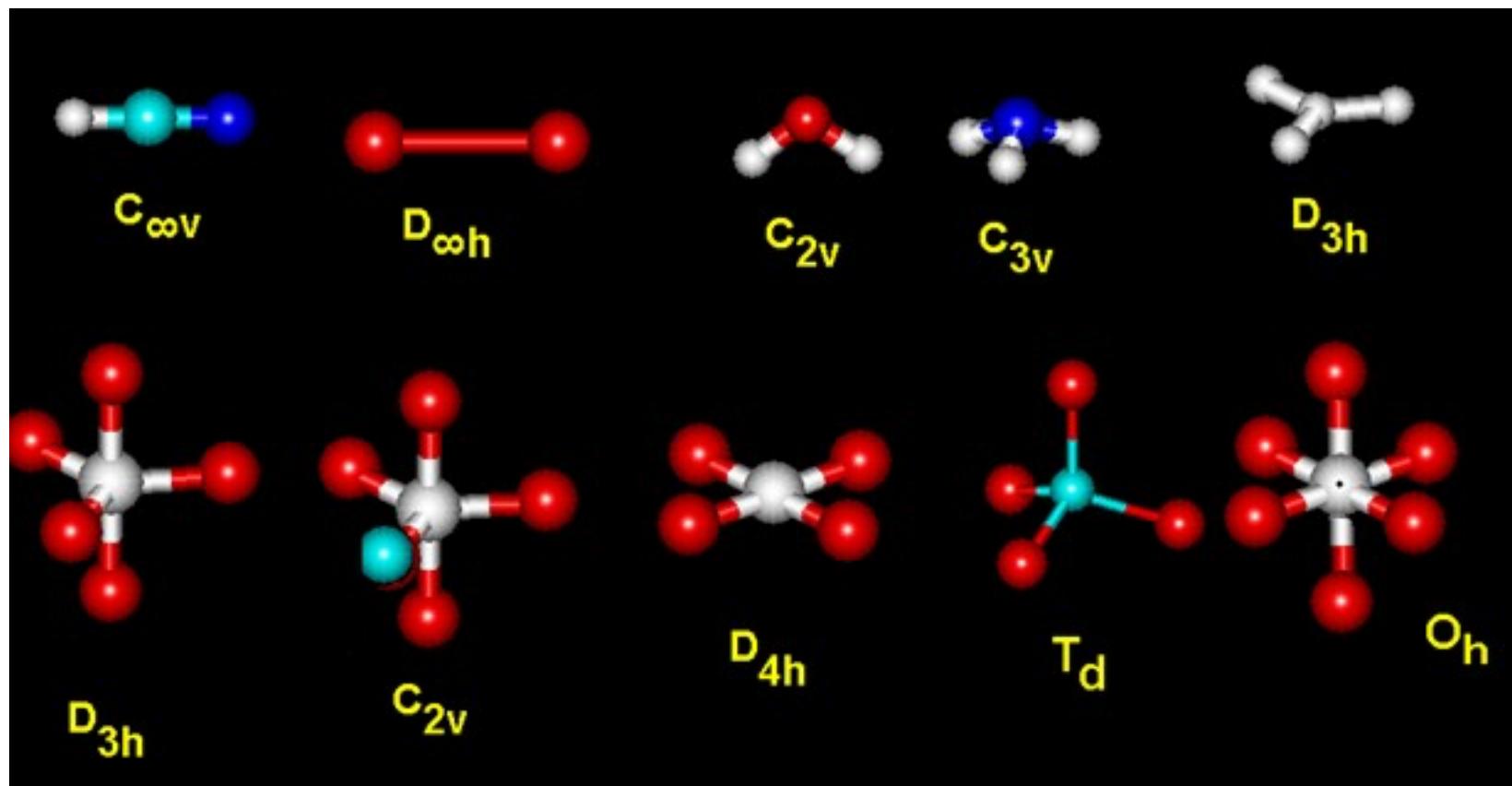
- **E**, the identity
- **C_n**, a rotation by an angle $2\pi/n$;
- **σ** reflection in a plane, classified as
 - **σ_h**, reflection through a plane perpendicular to the axis of highest rotation symmetry, called *principal axis*
 - **σ_v**, reflection through a plane to which the principal axis belongs
 - **σ_d**, reflection through a plane to which the principal axis belongs, and bisecting the angle between the two-fold axes perpendicular to the principal axis.
- **S_n = σ_h ⊗ C_n**, improper rotation of an angle $2\pi/n$
- **I = S₂**, the inversion.



Quanty : http://www.quanty.org/physics_chemistry/point_groups

Symmetries : some point groups

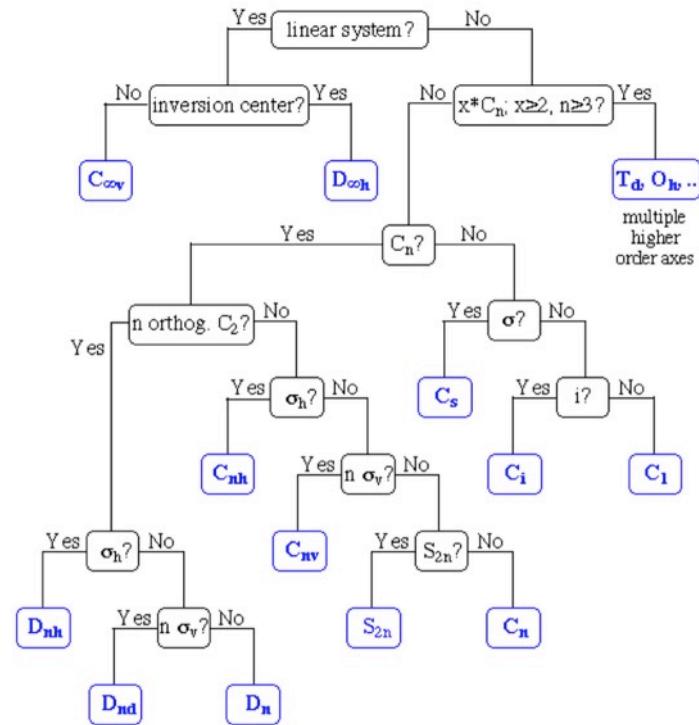
- ◆ A group of symmetry is an ensemble of symmetry operations (group theory)
- ◆ Each group is labeled (in Schönfliess notation)



- ◆ Each group has a table of representations

Symmetry tree

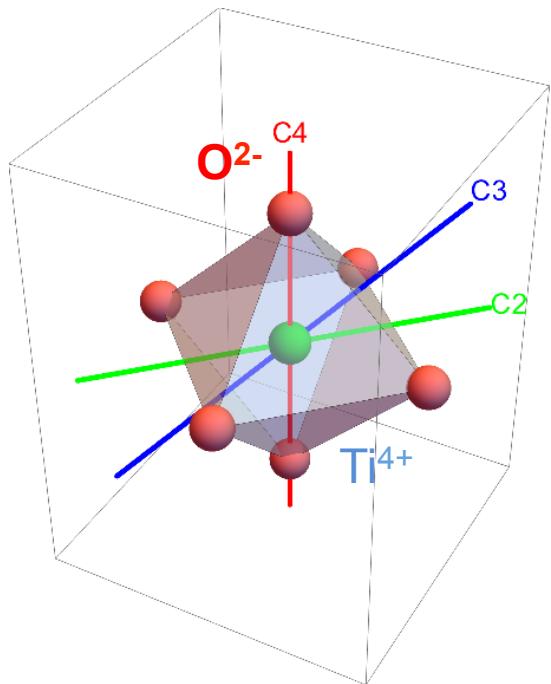
Symmetry - "Tree"



Quanty :

[http://www.quanty.org/physics_chemistry/
point_groups#a_flow_diagram_to_determine_the_point_group](http://www.quanty.org/physics_chemistry/point_groups#a_flow_diagram_to_determine_the_point_group)

Octahedral symmetry: O_h



Group : O_h

Character table

	Symmetry elements										Basis functions	
	E	$8C_3$	$6C_2$	$6C_4$	$3C_2 = (C_4)^2$	i	$6S_4$	$8S_6$	$3\sigma_h$	$6\sigma_d$	linear, rotations	quadratic
A_{1g}	1	1	1	1		1	1	1	1	1		$x^2+y^2+z^2$
A_{2g}	1	1	-1	-1		1	1	-1	1	1		
E_g	2	-1	0	0		2	2	0	-1	2	0	$(2z^2-x^2-y^2, x^2-y^2)$
T_{1g}	3	0	-1	1		-1	3	1	0	-1	-1	(R_x, R_y, R_z)
T_{2g}	3	0	1	-1		-1	3	-1	0	-1	1	(xz, yz, xy)
A_{1u}	1	1	1	1		1	-1	-1	-1	-1		
A_{2u}	1	1	-1	-1		1	-1	1	-1	-1		
E_u	2	-1	0	0		2	-2	0	1	-2	0	
T_{1u}	3	0	-1	1		-1	-3	-1	0	1	1	(x, y, z)
T_{2u}	3	0	1	-1		-1	-3	1	0	1	-1	



Irreducible representation :
labeled the symmetry properties of a state in the group

Point groups: Quanty website

All you need is in the web page Quanty.org

http://quanty.org/physics_chemistry/point_groups

The screenshot shows a Wikipedia-style page for "Point groups". The sidebar on the left contains a navigation menu with sections like "Navigation", "Quanty Documentation", "Physics and Chemistry" (which is highlighted with a blue box and has a blue arrow pointing to it), "Register", "Download area", "Copyright", "Authors", "How to cite", "User publications", "Questions and answers", "Forum", "Workshops", "Calendar", "Heidelberg", "Script versions", "Quanty.nb", "Quanty.m", "Quanty.py", "External programs", "CTM4", "Crispy", "DFT and Hartree-Fock", "Privacy", "Privacy policy", "Print/export", and "Printable version". The main content area has tabs for "Article" (selected) and "Talk". It features a "Table of Contents" box on the right with sections for Point groups, Different orientations, Symmetry operations, Irreducible representations, A flow diagram to determine the point group, Acknowledgements, Table of several point groups, and Table of contents. The central part of the page displays a table of point groups categorized by their symmetry elements:

Nonaxial groups	$C_1 - 1$	$C_s - m$	$C_i - \bar{1}$				
C_n groups	$C_2 - 2$	$C_3 - 3$	$C_4 - 4$	$C_5 - 5$	$C_6 - 6$	$C_7 - 7$	$C_8 - 8$
D_n groups	$D_2 - 222$	$D_3 - 32$	$D_4 - 422$	$D_5 - 52$	$D_6 - 622$	$D_7 - 722$	$D_8 - 822$
C_{nv} groups	$C_{2v} - mm2$	$C_{3v} - 3m$	$C_{4v} - 4mm$	$C_{5v} - 5m$	$C_{6v} - 6mm$	$C_{7v} - 7m$	$C_{8v} - 8mm$
C_{nh} groups	$C_{2h} - 2/m$	$C_{3h} - \bar{6}$	$C_{4h} - 4/m$	$C_{5h} - \bar{10}$	$C_{6h} - 6/m$		
D_{nh} groups	$D_{2h} - mmm$	$D_{3h} - \bar{6}m2$	$D_{4h} - 4/mmm$	$D_{5h} - \bar{10}m2$	$D_{6h} - 6/mmm$	$D_{7h} - \bar{14}m2$	$D_{8h} - 8/mmm$
D_{nd} groups	$D_{2d} - \bar{4}2m$	$D_{3d} - \bar{3}m$	$D_{4d} - \bar{8}2m$	$D_{5d} - \bar{5}m$	$D_{6d} - \bar{12}2m$	$D_{7d} - \bar{7}m$	$D_{8d} - \bar{16}2m$
S_n groups	$S_2 - \bar{2}$	$S_4 - \bar{4}$	$S_6 - \bar{6}$	$S_8 - \bar{8}$	$S_{10} - \bar{10}$	$S_{12} - \bar{12}$	
Cubic groups	$T - 23$	$T_h - m\bar{3}$	$T_d - \bar{4}3m$	$O - 432$	$O_h - m\bar{3}m$	$I - \bar{5}32$	$I_h - \bar{5}\bar{3}m$
Linear groups	C_{co}	D_{coh}					

The following set of pages list properties of the different point groups and their irreducible representations. The table above links to the main page of each of the different point group. For each group we list the character and product table. Often one needs to answer the question how a potential in a given point group looks like and what the eigen-states of that potential are. The eigen states can be grouped according to the irreducible representations of the group and for each of these representations one can give representing functions. The form of these functions and the potential however do depend on the orientation of the point group. We therefor list for each point group different orientations.

Different orientations

As we are interested in explicit representations we do need to specify the orientation of the symmetry operators. This results in several tables for the same point group but with different choices for the symmetry operations. For example the cubic O_h point group can be represented with the C_4 axes in the x , y and z direction, or with a C_3 axis in the z direction. We list several orientations of the different point-groups available.

Symmetry operations

We use the following notation for symmetry operations.

E = identity

C_n = n -fold rotation

S_n = n -fold rotation plus reflection through a plane perpendicular to the axis of rotation

Point groups: Quanty / Orientation

Oh

http://quanty.org/physics_chemistry/point_groups/oh

Navigation

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- Calendar
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Script versions

- Quanty.nb
- Quanty.m
- Quanty.py

External programs

- CTM4
- Crispy
- DFT and Hartree-Fock

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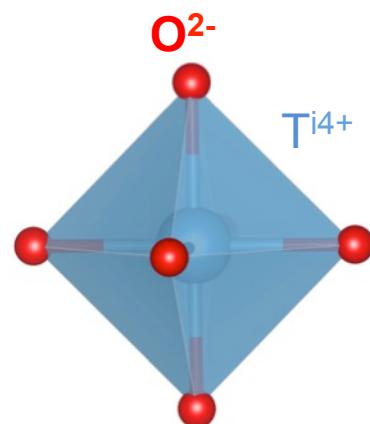
Orientation XYZ

Symmetry Operations

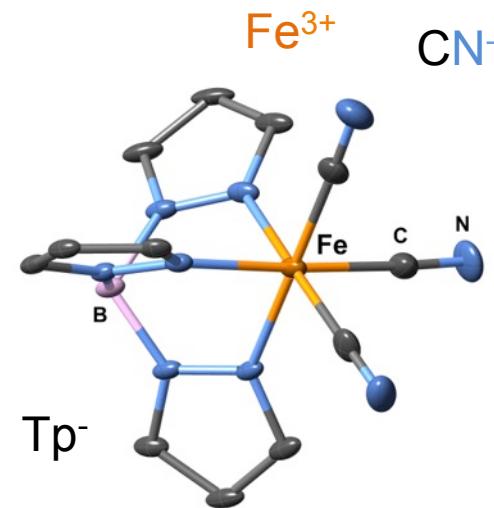
In the Oh Point Group, with orientation XYZ there are the following symmetry operations

Operator	Orientation
E	{0, 0, 0},
C_3	{1, 1, 1}, {1, 1, -1}, {1, -1, 1}, {-1, 1, 1}, {-1, -1, 1}, {-1, 1, -1}, {1, -1, 1}, {-1, -1, -1},
C_2	{1, 1, 0}, {1, -1, 0}, {1, 0, -1}, {1, 0, 1}, {0, 1, 1}, {0, 1, -1},
C_4	{0, 0, 1}, {0, 1, 0}, {1, 0, 0}, {0, 0, -1}, {0, -1, 0}, {-1, 0, 0},
C_2	{0, 0, 1}, {0, 1, 0}, {1, 0, 0},
i	{0, 0, 0},
S_4	{0, 0, 1}, {0, 1, 0}, {1, 0, 0}, {0, 0, -1}, {0, -1, 0}, {-1, 0, 0},
S_6	{1, 1, 1}, {1, 1, -1}, {1, -1, 1}, {-1, 1, 1}, {-1, -1, 1}, {-1, 1, -1}, {1, -1, 1}, {-1, -1, -1},
σ_h	{1, 0, 0}, {0, 1, 0}, {0, 0, 1},

Symmetries



Octahedral
Group : O_h



Approximative
Trigonal group :
 C_{3v}

II.2. Crystal field Hamiltonian

Crystal field Hamiltonian

CF potential : expanded on the normalized spherical harmonics

$$C_{k,m}(\theta, \phi) = \left(\frac{4\pi}{2k+1} \right)^{1/2} Y_{k,m}(\theta, \phi)$$

$$H_{CF} = \sum_{k=0}^{\infty} \sum_{m=-k}^{k} A_{k,m} C_k^m(\theta, \phi)$$

$A_{k,m}$ are the crystal field parameters

Sum over k infinity ?

→ CF matrix element

CF matrix element

The crystal field potential

$$H_{CF} = \sum_{k=0}^{\infty} \sum_{m=-k}^k A_{k,m} C_k^m(\theta, \phi)$$

Basis function (one electron)

$$\phi_i(\mathbf{r}, \sigma) = \frac{1}{r} P_{n_i \ell_i}(r) Y_{\ell, m}(\theta, \phi) \chi_{m_{s_i}}(\sigma)$$

Radial Angular Spherical harmonics spin

For an electron in orbital ℓ

$$\begin{aligned} \langle \phi_i | H_{CF} | \phi_j \rangle &\propto \langle Y_{\ell, m_i} | C_{k, m} | Y_{\ell, m_j} \rangle \\ &\propto \begin{pmatrix} \ell & k & \ell \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \ell & k & \ell \\ -m_i & m & m_j \end{pmatrix} \end{aligned}$$

$\neq 0$ if k even ($2\ell + k$ even), $0 \leq k \leq 2\ell$

Crystal field Hamiltonian

$$H_{CF} = \sum_{\substack{k=0, \\ k \text{ even}}}^{2\ell} \sum_{m=-k}^k A_{k,m} C_{k,m}(\theta, \phi)$$

\downarrow

$$A_{k,m} = (-1)^m A_{k,-m}^*$$

Hermitian Hamiltonian

Exemples

3d ion $\ell = 2$

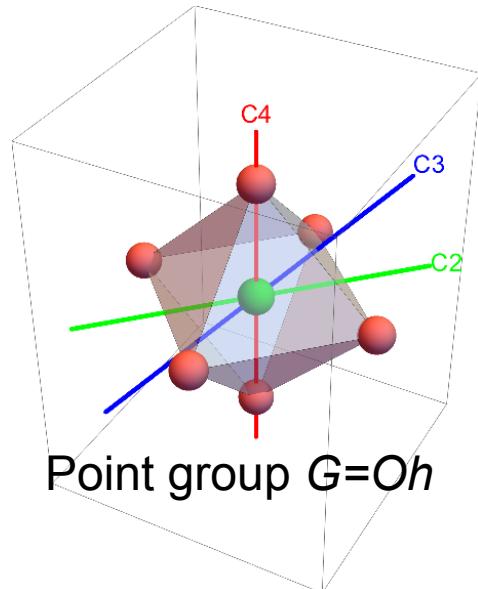
$k = 0, 2, 4$

4f ion $\ell = 3$

$k = 0, 2, 4, 6$

Crystal field and symmetry

$$H_{CF} = \sum_{\substack{k=0, \\ k \text{ even}}}^{2\ell} \sum_{m=-k}^k A_{k,m} C_{k,m}(\theta, \phi)$$



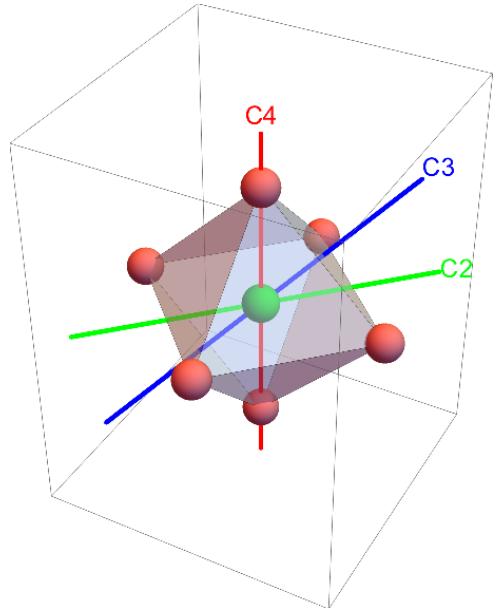
CF Hamiltonian belongs to G
invariant under all symmetry operation of $G(O_i)$

$$O_i H_{CF} = H_{CF}$$

Some $A_{km} = 0$ due to symmetry

Crystal field Hamiltonian and symmetry

Example : O_h symmetry



$$O_i H_{CF} = H_{CF}$$

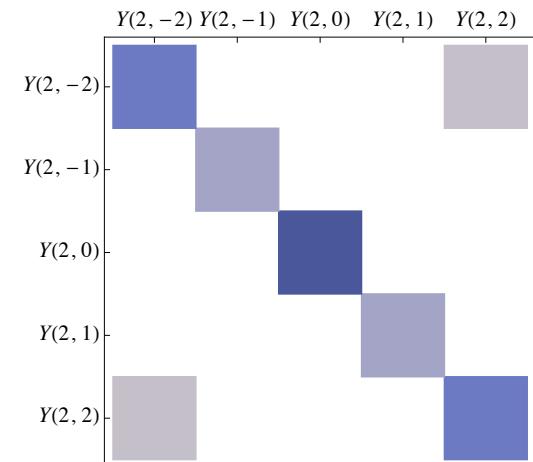
	E	$8C_3$	$6C_2$	$6C_4$	$3C_2 = (C_4)^2$	i	$6S_4$	$8S_6$	$3\sigma_h$	$6\sigma_d$
A_{1g}	1	1	1	1		1	1	1	1	1
A_{2g}	1	1	-1	-1		1	1	-1	1	-1
E_g	2	-1	0	0		2	2	0	-1	2
T_{1g}	3	0	-1	1		-1	3	1	0	-1
T_{2g}	3	0	1	-1		-1	3	-1	0	-1
A_{1u}	1	1	1	1		1	-1	-1	-1	-1
A_{2u}	1	1	-1	-1		1	-1	1	-1	1
E_u	2	-1	0	0		2	-2	0	1	-2
T_{1u}	3	0	-1	1		-1	-3	-1	0	1
T_{2u}	3	0	1	-1		-1	-3	1	0	1

$$H_{CF} \in A_{1g}$$

Only 2 non-zero A_{km} :
 A_{00}, A_{40}

$$H_{CF}^{O_h} = A_{4,0}C_4^0 + \sqrt{\frac{5}{14}}A_{4,0}(C_4^{-4} + C_4^4)$$

CF matrix



Notations for crystal field parameters

Vary from an author to the other

$$H_{CF} = \sum_{\substack{k=0, \\ k \text{ even}}}^{2\ell} \sum_{m=-k}^k A_{k,m} C_{k,m}(\theta, \phi)$$

(Haverkort/Quanty)

$$H_{CF}(\mathbf{r}) = \sum_{k=0}^{\infty} \sum_{m=-k}^k r^k A_{k,m} C_{k,m}(\theta, \phi) = \sum_{k=0}^{\infty} \sum_{m=-k}^k B_{k,m} C_{k,m}(\theta, \phi)$$

Parameters used mostly by chemist
Symmetry-dependent
(Balhausen, König, Kremer,...)

10Dq for Oh
10Dq, Ds, Dt for D_{4h}
10Dq, Dσ, Dτ for D_{3d} or C_{3v}

Parameters used in Quanty :
Orbital energies + off-diagonal elements
Symmetry-dependent

E_{eg}, E_{t2g} for Oh
E_{a1g}, E_{eg}, E_{b1g}, E_{b2g} for D_{4h}

Crystal field

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- 3. Crystal field splitting of one 3d electron**
3. Multi-electrons ion (3d electrons)
 1. The spherical ion :
 2. Crystal field
 3. Core hole spectroscopy : examples
4. *Crystal field and magnetism*
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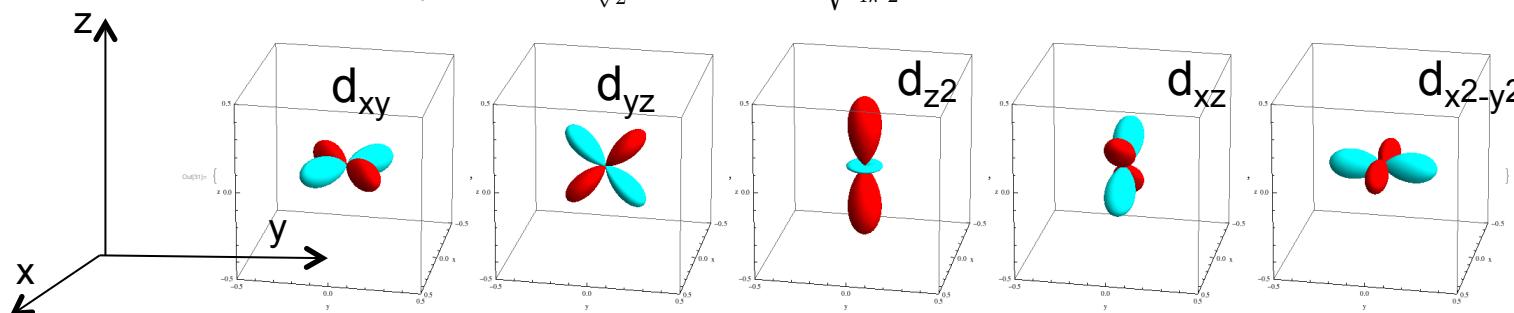
3d atomic orbitals

- Basis of atomic orbitals $\phi_i(\mathbf{r}, \sigma) = \frac{1}{r} P_{n_i \ell_i}(r) Y_{\ell, m}(\theta, \phi) \chi_{m_{s_i}}(\sigma)$
- 3d shell : n=3, $\ell=2$ $2\ell+1=5$ basis functions ($-2 \leq m_\ell \leq 2$)

$Y_{2,m}(\theta, \phi)$ with $m = -2, -1, 0, 1, 2$

- **3d orbitals** = real functions, linear combination of $Y_{2,m}(\theta, \phi)$

$$\begin{aligned}
 d_{xy} &= y_{2-2} = \frac{i}{\sqrt{2}}(Y_2^2 - Y_{-2}^2) = \sqrt{\frac{15}{4\pi}} \quad xy/r^2 \\
 d_{yz} &= y_{2-1} = \frac{i}{\sqrt{2}}(Y_1^2 + Y_{-1}^2) = \sqrt{\frac{15}{4\pi}} \quad yz/r^2 \\
 d_{3z^2-r^2} &= y_{20} = Y_2^0 = \sqrt{\frac{15}{4\pi}} \frac{1}{2\sqrt{3}} (3z^2 - r^2)/r^2 \\
 d_{xz} &= y_{21} = \frac{1}{\sqrt{2}}(Y_1^2 - Y_{-1}^2) = \sqrt{\frac{15}{4\pi}} \quad xz/r^2 \\
 d_{x^2-y^2} &= y_{22} = \frac{1}{\sqrt{2}}(Y_2^2 + Y_{-2}^2) = \sqrt{\frac{15}{4\pi}} \frac{1}{2} (x^2 - y^2)/r^2
 \end{aligned}$$



Crystal field splitting for O_h

$$H_{CF}^{O_h} = A_{4,0}C_{4,0} + \sqrt{\frac{5}{14}}A_{4,0}(C_{4,-4} + C_{4,4}) = 21D_qC_{4,0} + 21\sqrt{\frac{5}{14}}D_q(C_{4,-4} + C_{4,4})$$

1) CF matrix in $\{Y_{2,m}\}$ basis

$$\begin{array}{ccccc}
 |\ell m\rangle \equiv |2-2\rangle & |2-1\rangle & |2\ 0\rangle & |2\ 1\rangle & |2\ 2\rangle \\
 \left(\begin{array}{ccccc}
 A_{0,0} + \frac{1}{21}A_{4,0} & 0 & 0 & 0 & \frac{5}{21}A_{4,0} \\
 0 & A_{0,0} - \frac{4}{21}A_{4,0} & 0 & 0 & 0 \\
 0 & 0 & A_{0,0} + \frac{2}{7}A_{4,0} & 0 & 0 \\
 0 & 0 & 0 & A_{0,0} - \frac{4}{21}A_{4,0} & 0 \\
 \frac{5}{21}A_{4,0} & 0 & 0 & 0 & A_{0,0} + \frac{1}{21}A_{4,0}
 \end{array} \right)
 \end{array}$$

2) Diagonalization

$$\begin{array}{ccccc}
 d_{x^2-y^2} & d_{z^2} & d_{yz} & d_{xz} & d_{xy} \\
 \left(\begin{array}{ccccc}
 A_{0,0} + \frac{2}{7}A_{4,0} & 0 & 0 & 0 & 0 \\
 0 & A_{0,0} + \frac{2}{7}A_{4,0} & 0 & 0 & 0 \\
 0 & 0 & A_{0,0} - \frac{4}{21}A_{4,0} & 0 & 0 \\
 0 & 0 & 0 & A_{0,0} - \frac{4}{21}A_{4,0} & 0 \\
 0 & 0 & 0 & 0 & A_{0,0} - \frac{4}{21}A_{4,0}
 \end{array} \right)
 \end{array}$$

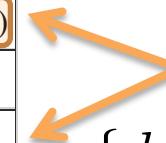
→ O_h crystal field splits the d orbitals in two groups

$\{d_{x^2-y^2}, d_{z^2}\}$ and $\{d_{yz}, d_{xz}, d_{xy}\}$

O_h Crystal field splitting of 3d electron Use of group symmetry

Character table for O_h point group

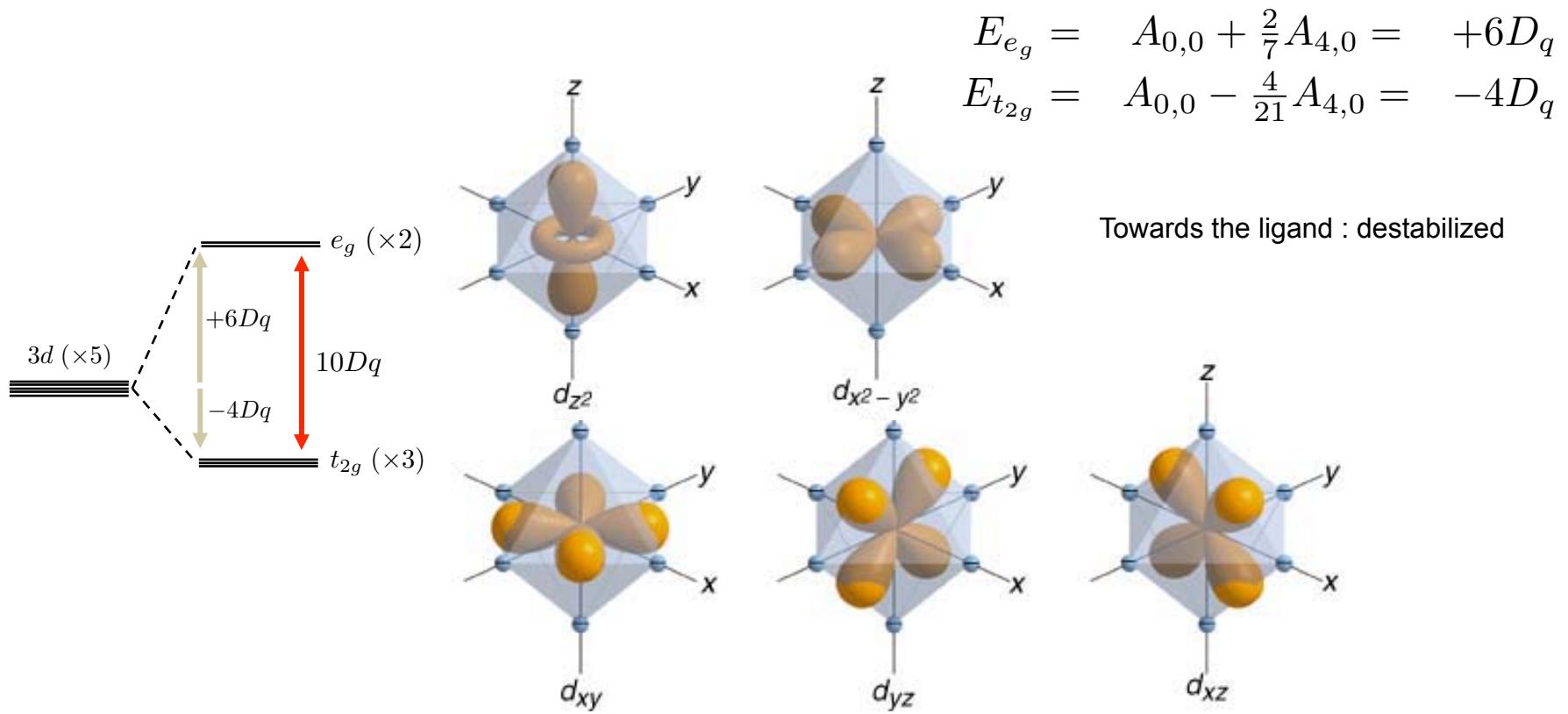
	E	8C ₃	6C ₂	6C ₄	3C ₂ = (C ₄) ²	i	6S ₄	8S ₆	3σ _h	6σ _d	linear, rotations	quadratic
A _{1g}	1	1	1	1		1	1	1	1	1		x ² +y ² +z ²
A _{2g}	1	1	-1	-1		1	1	-1	1	1	-1	
E _g	2	-1	0	0		2	2	0	-1	2	0	(2z ² -x ² -y ² , x ² -y ²)
T _{1g}	3	0	-1	1		-1	3	1	0	-1	-1	(R _x , R _y , R _z)
T _{2g}	3	0	1	-1		-1	3	-1	0	-1	1	(xz, yz, xy)
A _{1u}	1	1	1	1		1	-1	-1	-1	-1	-1	
A _{2u}	1	1	-1	-1		1	-1	1	-1	-1	1	
E _u	2	-1	0	0		2	-2	0	1	-2	0	
T _{1u}	3	0	-1	1		-1	-3	-1	0	1	1	(x, y, z)
T _{2u}	3	0	1	-1		-1	-3	1	0	1	-1	


 $\{d_{x^2-y^2}, d_{z^2}\} \in e_g$
 $\{d_{yz}, d_{xz}, d_{xy}\} \in t_{2g}$

- From the O_h group properties, one can guess the splitting of the d orbitals
- d orbitals are called e_g and t_{2g} *

(*)N.B. : lower case letter (e_g and t_{2g}) for one electron

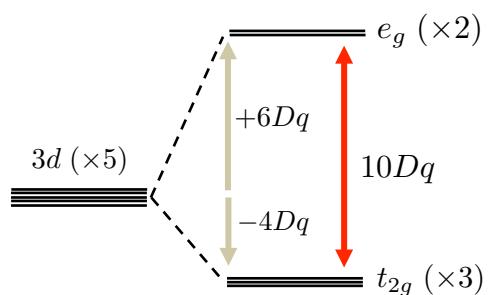
O_h Crystal field splitting of 3d electron



$10Dq \equiv$ Crystal field strength

(parameter)

O_h Crystal field splitting of 3d electron Parameters



$$E_{e_g} = A_{0,0} + \frac{2}{7}A_{4,0} = +6D_q$$

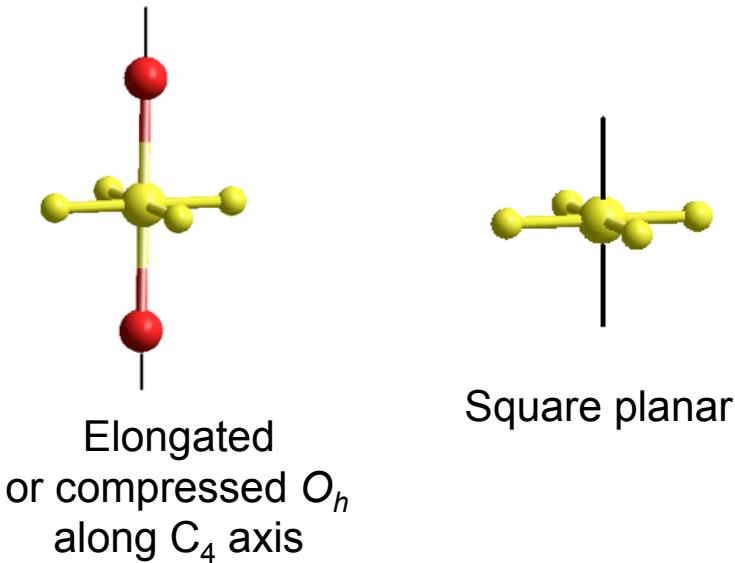
$$E_{t_{2g}} = A_{0,0} - \frac{4}{21}A_{4,0} = -4D_q$$

Quantity : A_{k,m} defined in function the orbital energies

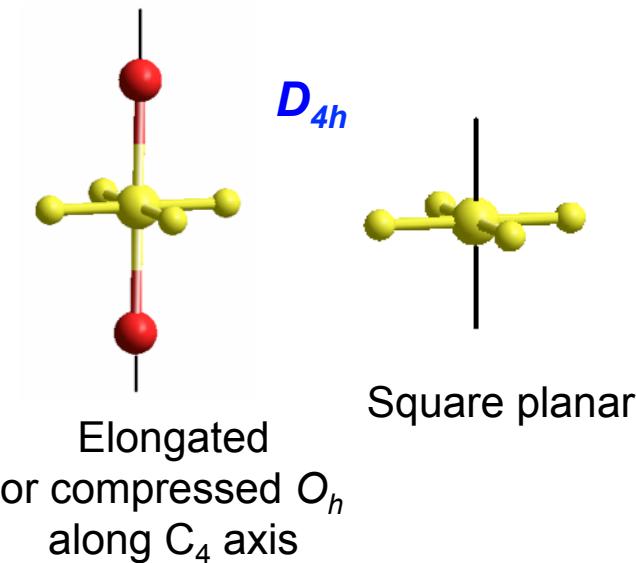
$$A_{k,m} = \begin{cases} \frac{2}{5}E_{e_g} + \frac{3}{5}E_{t_{2g}} & k = 0 \quad m = 0 \\ \frac{21}{10}(E_{e_g} - E_{t_{2g}}) & k = 4 \quad m = 0 \\ \frac{21}{10}\sqrt{\frac{5}{14}}(E_{e_g} - E_{t_{2g}}) & k = 4 \quad m = \pm 4 \\ 0 & \text{True} \end{cases}$$

Tetragonal (D_{4h}) crystal field

D_{4h}



Tetragonal (D_{4h}) crystal field



Character table for D_{4h} point group

	E	$2C_4$ (z)	C_2	$2C'_2$	$2C''_2$	i	$2S_4$	σ_h	$2\sigma_v$	$2\sigma_d$	linears, rotations	quadratic
A_{1g}	1	1	1	1	1	1	1	1	1	1		x^2+y^2, z^2
A_{2g}	1	1	1	-1	-1	1	1	1	-1	-1	R_z	
B_{1g}	1	-1	1	1	-1	1	-1	1	1	-1		x^2-y^2
B_{2g}	1	-1	1	-1	1	1	-1	1	-1	1		xy
E_g	2	0	-2	0	0	2	0	-2	0	0	(R_x, R_y)	(xz, yz)
A_{1u}	1	1	1	1	1	-1	-1	-1	-1	-1		
A_{2u}	1	1	1	-1	-1	-1	-1	-1	1	1	z	
B_{1u}	1	-1	1	1	-1	-1	1	-1	-1	1		
B_{2u}	1	-1	1	1	1	-1	1	-1	1	-1		
E_u	2	0	-2	0	0	-2	0	2	0	0		(x, y)

From D_{4h} table, one predicts the 3d splitting in **4 groups**

$$\{d_{z^2}\} \in a_{1g}$$

$$\{d_{x^2-y^2}\} \in b_{1g}$$

$$\{d_{xy}\} \in b_{2g}$$

$$\{d_{xz}, d_{yz}\} \in e_g$$

Tetragonal (D_{4h}) crystal field

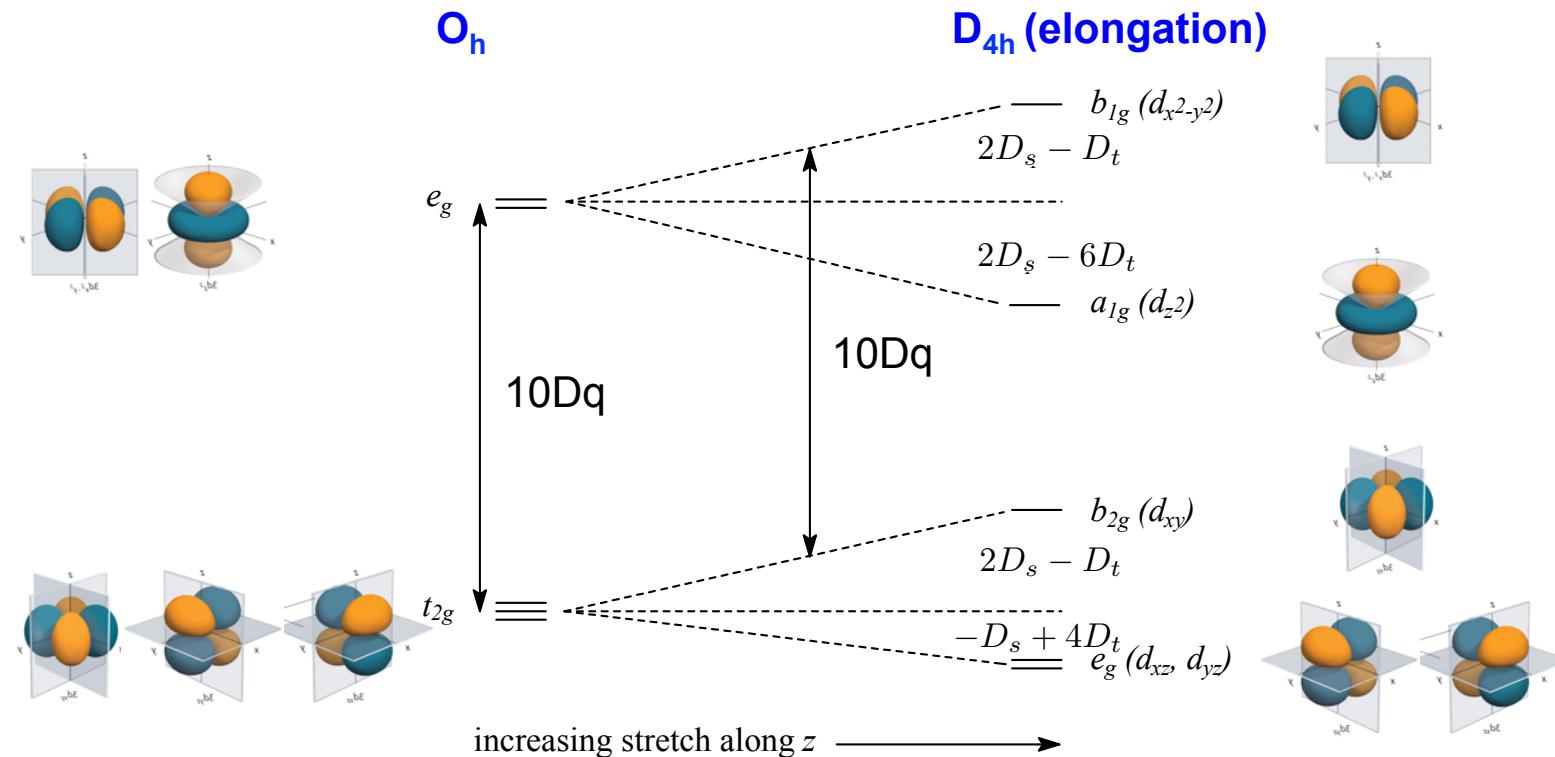
$$H_{CF}^{D_{4h}} = A_{0,0}C_{0,0} + A_{2,0}C_{2,0} + A_{4,0}C_{4,0} + A_{4,4}(C_{4,-4} + C_{4,4})$$

CF matrix diagonal in the $\{d\}$ basis

$$\begin{pmatrix} d_{x^2-y^2} & \color{red}d_{z^2} & \color{blue}d_{yz} & \color{blue}d_{xz} & \color{green}d_{xy} \\ E_{a_{1g}} & 0 & 0 & 0 & 0 \\ 0 & \color{red}E_{b_{1g}} & 0 & 0 & 0 \\ 0 & 0 & \color{blue}E_{e_g} & 0 & 0 \\ 0 & 0 & 0 & \color{blue}E_{e_g} & 0 \\ 0 & 0 & 0 & 0 & \color{green}E_{b_{2g}} \end{pmatrix} \quad \text{4 energy levels}$$

Tetragonal (D_{4h}) crystal field splitting of 3d electron

$$H_{CF}^{D_{4h}} = A_{0,0}C_{0,0} + A_{2,0}C_{2,0} + A_{4,0}C_{4,0} + A_{4,4}(C_{4,-4} + C_{4,4})$$



(*)The relation with Dq, Ds, Dt in König&Kremer « Ligand field. Energy diagram »

Crystal field Hamiltonian in Quanty

$$H_{CF}^{O_h} = A_{0,0}C_{0,0} + A_{4,0}C_{4,0} + \sqrt{\frac{5}{14}}A_{4,0}(C_{4,-4} + C_{4,4}) = 21D_qC_{4,0} + 21\sqrt{\frac{5}{14}}D_q(C_{4,-4} + C_{4,4})$$

1) Pre-defined CF potential

```
Akm = PotentialExpandedOnYlm("Oh", 2, {0.6, -0.4});  
OpptenDq = NewOperator("CF", NFermion, dIndexUp, dIndexDn, Akm);
```

The Oh potential is defined by:

$$A_{k,m} = \begin{cases} \frac{2}{5}E_{e_g} + \frac{3}{5}E_{t_{2g}} & k = 0 \quad m = 0 \\ \frac{21}{10}(E_{e_g} - E_{t_{2g}}) & k = 4 \quad m = 0 \\ \frac{21}{10}\sqrt{\frac{5}{14}}(E_{e_g} - E_{t_{2g}}) & k = 4 \quad m = \pm 4 \\ 0 & \text{True} \end{cases}$$

1) User made CF potential: Akm= $\{(k_1, m_1, A_{k_1, m_1}), (k_2, m_2, A_{k_2, m_2}), \dots\}$

```
Akm = {{4, 0, 21/10}, {4, -4, 21/10sqrt(5/14)}, {4, 4, 21/10sqrt(5/14)}};  
OpptenDq = NewOperator("CF", NFermion, dIndexUp, dIndexDn, Akm);
```

The parameter is 10Dq and the Hamiltonian writes $H_{CF} = 10Dq * OpptenDq$

To lower symmetries

O_h	1 parameter (10Dq)	2 energies
D_{4h}	3 parameters (10Dq,Ds,Dt)	4 energies
D_{3d}	3 parameters (10Dq,D σ ,D τ)	3 energies
C_{2v}	9 parameters	5 energies (5 orbitals) 

Crystal field

1. Introduction / Definition
2. One electron (3d electron)

3. Multi-electrons ion (3dⁿ)

1. The spherical ion :
 1. Configuration
 2. Hamiltonian
 3. $|(L,S),J\rangle$ basis function
 4. Spectroscopic terms
2. Crystal field
 1. Basis functions, Hamiltonian and matrix element
 2. Energy diagram (Tanabe-Sugano)
3. Core hole spectroscopy : examples
 1. Spin cross-over
 2. The limit of the crystal field model : towards ligand field
4. *Crystal field and magnetism*

Multi-electrons ions $3d^n$

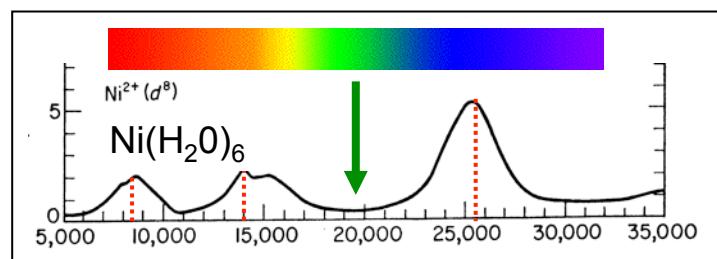
A lot of physical-chemical properties are due to the crystal field

For example

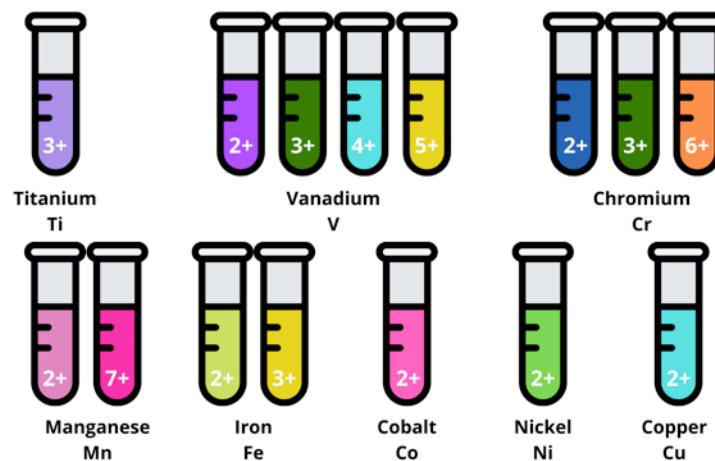
- Color
- Magnetic properties



- Atomic number (Z)
- Number of electron (n) $3d^n$
- Crystal field strength ($\approx 10Dq$) (ligand)



Transition Metal Ion Colors in Aqueous Solution



UV-visible Absorption spectroscopy

Multi-electrons ions : configuration

- | | | |
|------------------------------------|---|-------------------------|
| ➤ Electronic configuration: | complete | simplified (open shell) |
| 3d transition metal ions | $1s^2 2s^2 2p^6 3d^n$ | $3d^n$ |
| 4f rare earth ions | $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4f^n$ | $4f^n$ |

- ### ➤ Filling of the 3d orbitals:

$2\ell+1=5$ $|\ell, m_\ell\rangle$ angular functions
 $2s+1=5$ $|sm_s\rangle$ spin functions } 10 functions

$$\text{Nb states} = \binom{10}{n} = \frac{10!}{n!(10-n)!}$$

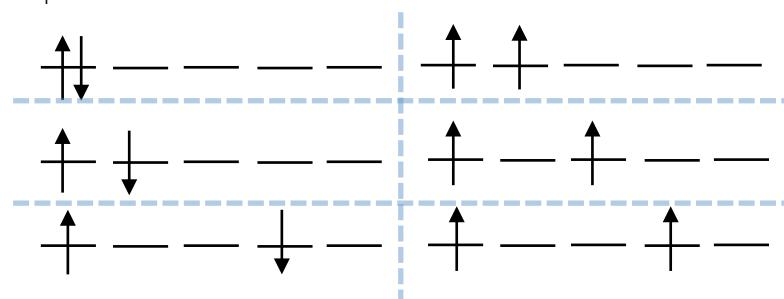
Example :

3d² ion (V³⁺, Cr⁴⁺):

degenerate states : $C_{10}^2 = 45$

Some possibilities:

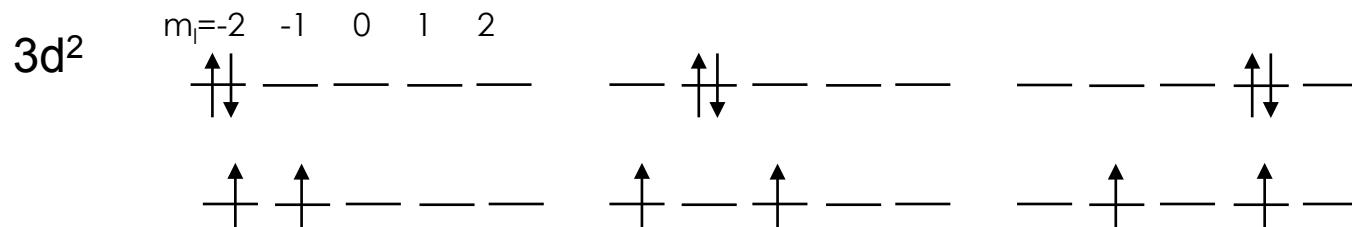
$$m_{\text{I}} = -2 \quad -1 \quad 0 \quad 1 \quad 2$$



Multiplets

n electron in 10 d orbitals

$$\text{Nb states} = \frac{10!}{n!(10-n)!}$$



Degeneracy => **Multiplets (doublet, triplet,...)**

→ Name of the model :

Crystal Field Multiplet (CFM)
Ligand Field Multiplet (LFM)

Hamiltonian of the ion with N electrons *Crystal field*

$$H_{\text{ion}} = H_{\text{cin}} + H_{e-n} + H_{e-e} + H_{s-o} + H_{\text{CF}}$$

{
 Free ion (spherical) Crystal field

$$H_{\text{cin}} = \sum_{i=1}^N -\frac{\hbar^2}{2m} \nabla_{r_i}^2 \quad \text{Total kinetic energy}$$

$$H_{e-n} = \sum_{i=1}^N -\frac{Ze^2}{4\pi\epsilon_0 r_i} \quad \text{Coulomb attraction nuclei-electrons}$$

$$H_{e-e} = \sum_{i < j=1}^N \frac{e^2}{4\pi\epsilon_0 r_{ij}} \quad \text{Electron-electron Coulomb repulsions}$$

$$H_{s-o} = \sum_{i=1}^N \xi_i(r_i) \mathbf{l}_i \cdot \mathbf{s}_i \quad \text{Spin-orbit coupling}$$

Multi-electrons ions configuration and symmetry

d² ion in O_h symmetry

e _g	— —	— —	↑ —	↓ —
t _{2g}	↑↓ — —	↑ +↓ —	↑ — —	↑ — —
Orbital degeneracy	3	6	3x2=6	3x2=6
Spin degeneracy	1 (S=0)	1 (S=0)	3 (S=1)	1 (S=0)

Hamiltonian of the ion with N electrons

Basis functions $|\Psi\rangle$

- **Multi-electron functions for N electrons (configuration)**

$$\Psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N)$$

- **Linear combination of Slater determinant built from the mono-electronic functions**

$$\phi_i(\mathbf{r}, \sigma) = \frac{1}{r} P_{n_i \ell_i}(r) Y_{\ell, m}(\theta, \phi) \chi_{m_{s_i}}(\sigma)$$

- **Anti-symmetric function to satisfy the Pauli principle**

Free ion with N electrons : basis functions

- Functions $|\alpha(L,S)JM_J\rangle = (L,S)J$ coupled functions

✓ For N electrons, $\hat{L} = \sum_{i=1}^N \hat{\ell}_i$  Sum is coupling

 $\hat{\ell}_i \neq \ell_i$

$\hat{\ell}_i$ is the orbit operator defined by

$$\hat{\ell}_{z,i} |\ell_i m_{\ell_i}\rangle = m_{\ell_i} |\ell_i m_{\ell_i}\rangle$$

$$\hat{\ell}_i^2 |\ell_i m_{\ell_i}\rangle = \ell_i(\ell_i + 1) |\ell_i m_{\ell_i}\rangle$$

$$-\ell_i \leq m_{\ell_i} \leq \ell_i$$

- ✓ Same definitions for \hat{S} and \hat{J}
- ✓ $\hat{J} = \hat{L} \oplus \hat{S}$ operator associated with the spin-orbit coupling
 $L - S \leq J \leq |L + S|$

Free $3d^n$ ion : basis functions, spectroscopic terms

No spin-orbit Basis functions $|\mathbf{L}, \mathbf{S}, M_L, M_S\rangle$

$H_{cin} + H_{e-n} + H_{e-e}$ commutes with $\hat{L}^2, \hat{L}_z, \hat{S}^2, \hat{S}_z$

Common eigenfunctions with \hat{L}^2, \hat{S}^2

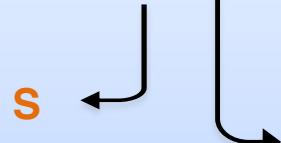
The energy of $|\alpha(L, S)\rangle$ does not depend on M_L, M_S

Degeneracy of $|\alpha(L, S)\rangle = (2L+1) (2S+1)$

Orbital degeneracy Spin degeneracy

Spectroscopic terms $|\mathbf{L}, \mathbf{S}, M_L, M_S\rangle$

$2S+1 \Gamma$



letter for L

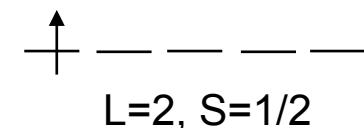
$\Gamma=S$ pour $L=0$

$\Gamma=P$ pour $L=1$

$\Gamma=D$ pour $L=2$

$\Gamma=F$ pour $L=3$

Ex: d^1



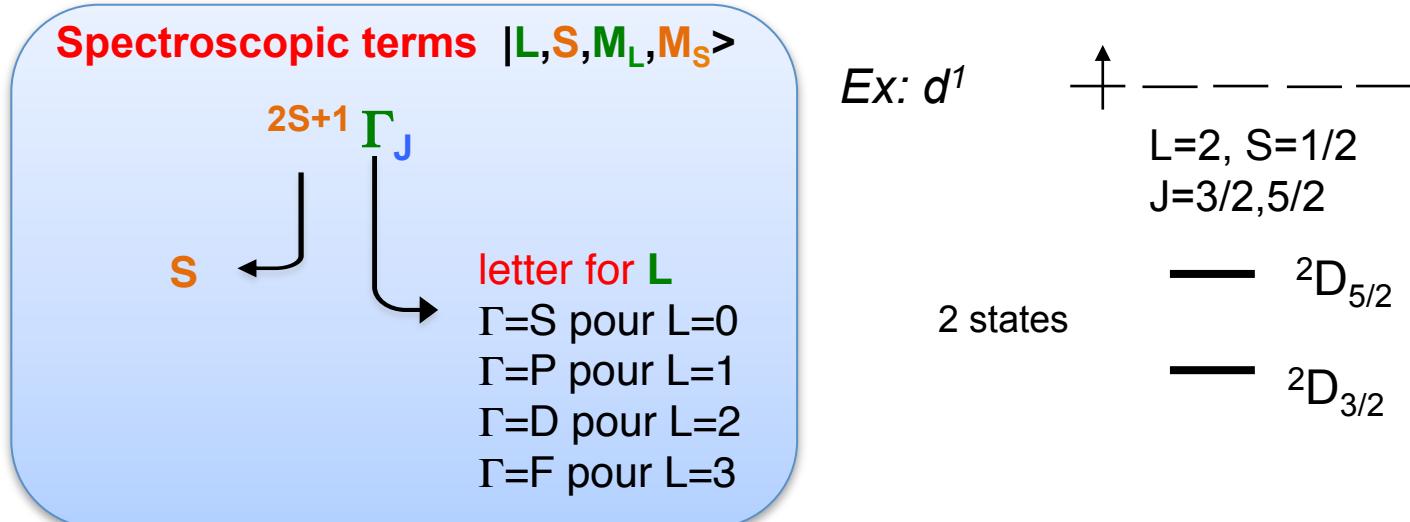
1 state

2D

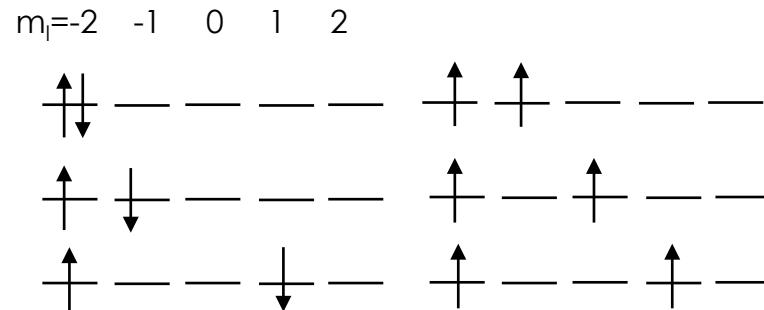
Free ion with N electrons : basis functions

Spin-orbit Basis functions $|(L,S)JM_J\rangle$

- Common eigenfunctions with $\hat{J}^2, \hat{J}_z, \hat{L}^2, \hat{S}^2$
 $H_{cin} + H_{e-n} + H_{e-e} + H_{s-0}$ commutes with $\hat{J}^2, \hat{J}_z, \hat{L}^2, \hat{S}^2$
 $(H_{cin} + H_{e-n} + H_{e-e})$ commutes with $\hat{L}^2, \hat{L}_z, \hat{S}^2, \hat{S}_z$)
- The energy of $|\alpha(L,S)JM_J\rangle$ does not depend on M_J
Degeneracy of $|\alpha(L,S)JM_J\rangle = 2J+1$



Free ion : $3d^2$ ion



$$\hat{L} = \hat{\ell} \otimes \hat{\ell} \quad 0 \leq L \leq 2\ell = 4$$

S, P, D, F, G terms

$$\hat{S} = \hat{s} \otimes \hat{s} \quad 0 \leq S \leq 2s = 1$$

Spin doublet or triplet

5 terms

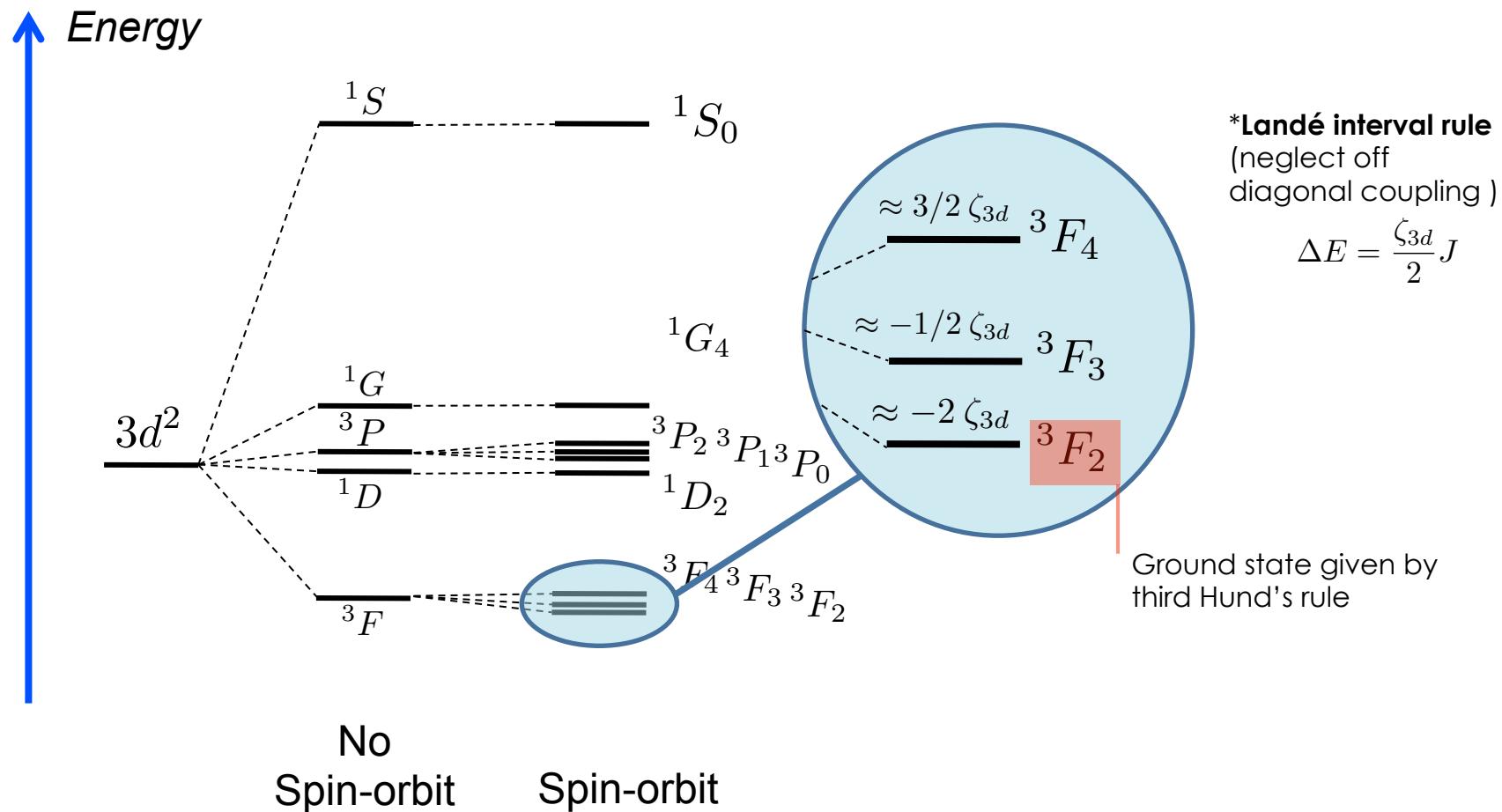
${}^1G, {}^3F, {}^1D, {}^3P, {}^1S$



Ground state (Hund's rules)

Free ion

Spectroscopic terms for $3d^2$ ion



Crystal field : basis function

➤ Basis functions

Spherical symmetry:

$$|\alpha(L, S) J M_J\rangle$$

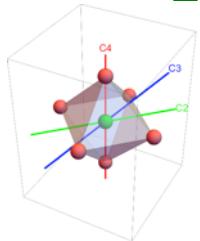
Local symmetry around the ion → point group G:

Basis functions of the representations of group G

$$|\alpha(L, S) J \Gamma \gamma\rangle$$

Γ irreducible representation of G

$\Gamma\gamma$ basis function of G



➤ Thole's code uses group theory and the $\{\Gamma\}$ basis (TTMULT, CTM4XAS)

➤ Quanty

- uses group theory only for crystal field potential building
- uses **spherical $\{Y_{l,m}\}$ basis** and not the $\{\Gamma\}$ basis.
- although Quanty is a many body code, the basis set is defined by one particle
- uses $\{\Gamma\}$ to label the states (spectroscopic terms)

Crystal field and group theory

Matrix elements

$$H_{CF} \in A_{1(g)}^*$$

(fully symmetric representation of group **G**)

* g only for centro-symmetric group



$$\langle (L_i, S_i) J_i \Gamma_i | H_{CF} | (L_j, S_j) J_j \Gamma_j \rangle \neq 0$$

if $\Gamma_i \otimes \Gamma_{H_{CF}} \otimes \Gamma_j \ni A_{1(g)}$

if $\Gamma_i \otimes A_{1(g)} \otimes \Gamma_j = \Gamma_i \otimes \Gamma_j \ni A_{1(g)}$

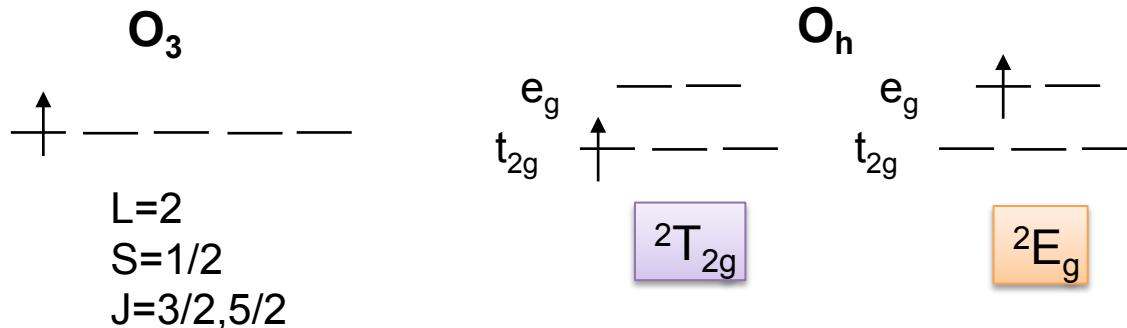


Crystal field mixes J states if $\Gamma_i \otimes \Gamma_j \ni A_{1(g)}$

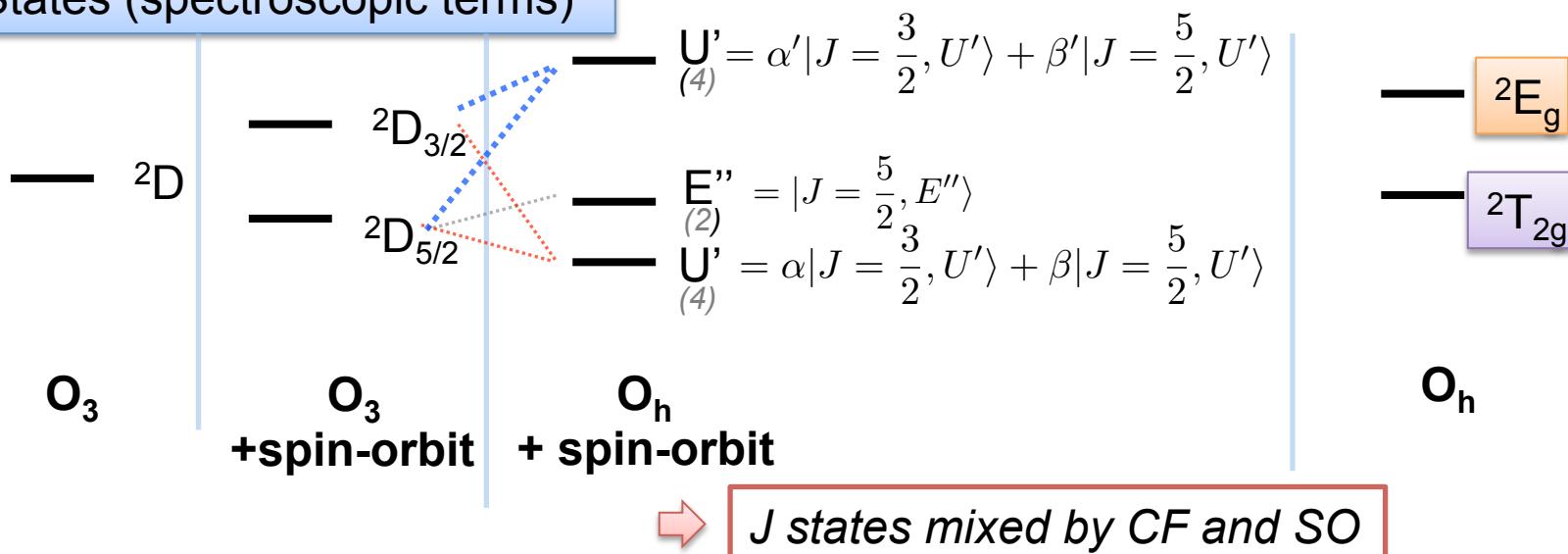
(group multiplication table)

Ex: d^1 (d^9) ion in O_h symmetry

Orbitals

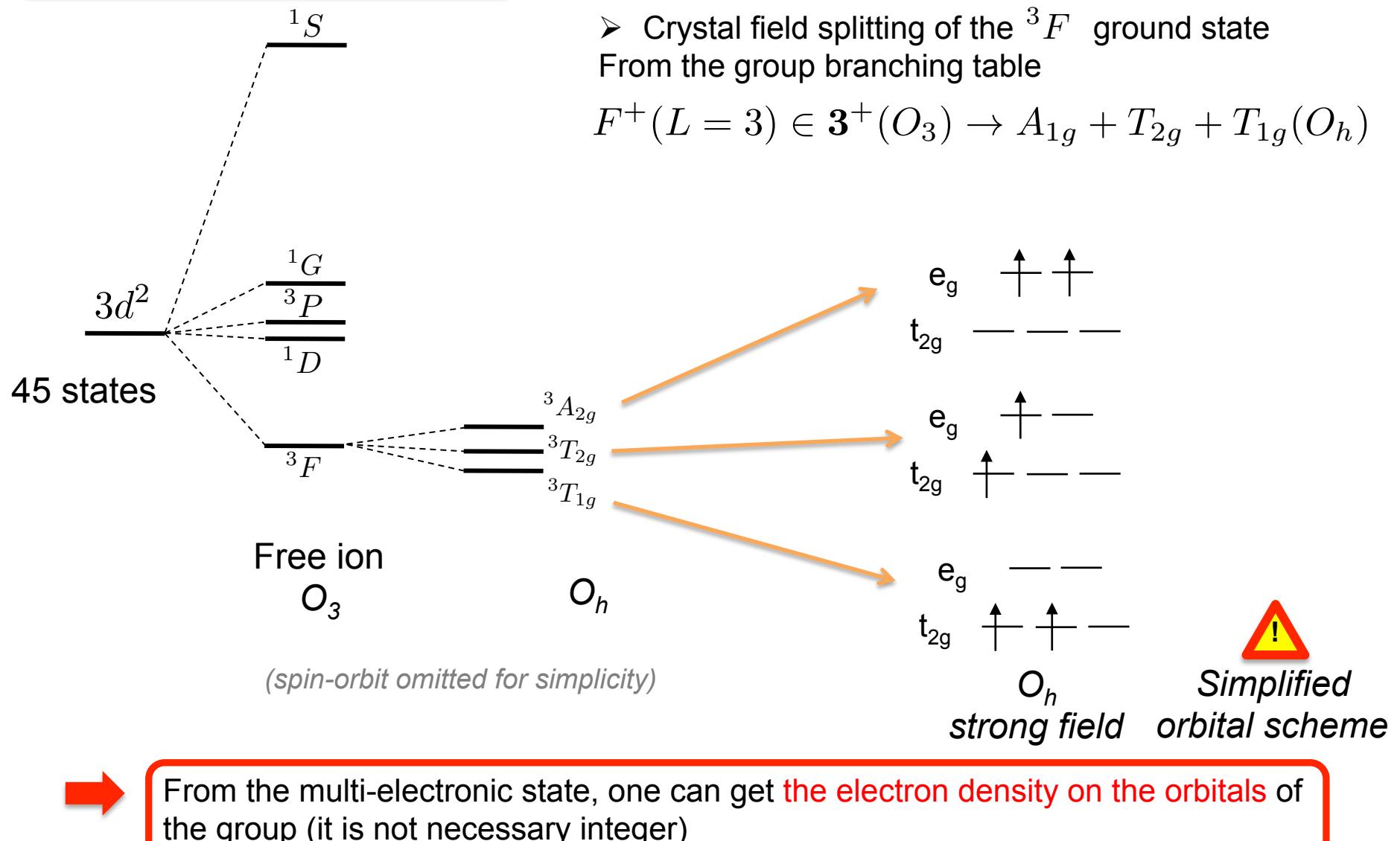


States (spectroscopic terms)



Crystal field and group theory

Example: d^2 ion in O_h symmetry

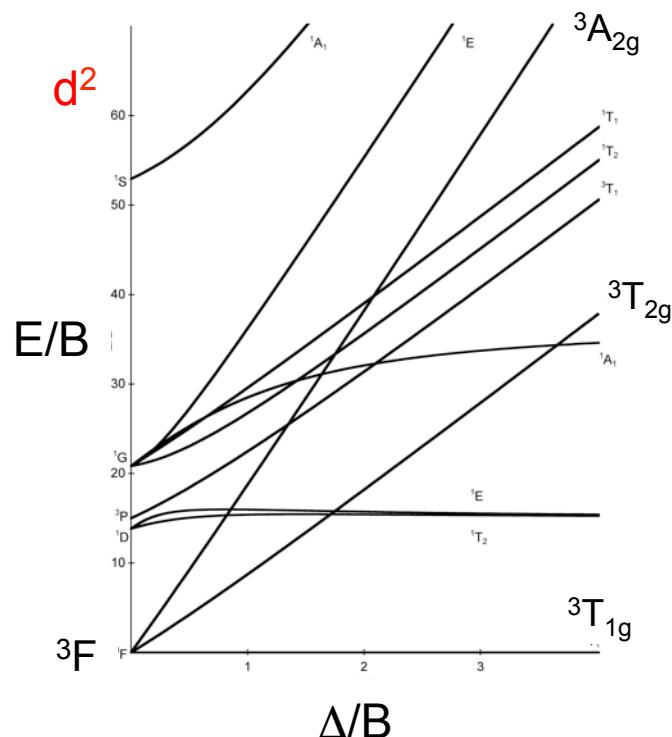


Crystal field : energy diagram

- Plot of the **energy** of spectroscopic terms ($^{2S+1}\Gamma$ or $^{2S+1}\Gamma_J$) as function of **crystal field parameter (10Dq,Ds,...)** (and B Racah parameter (*))

- **Tanabe-Sugano diagram (1954)**

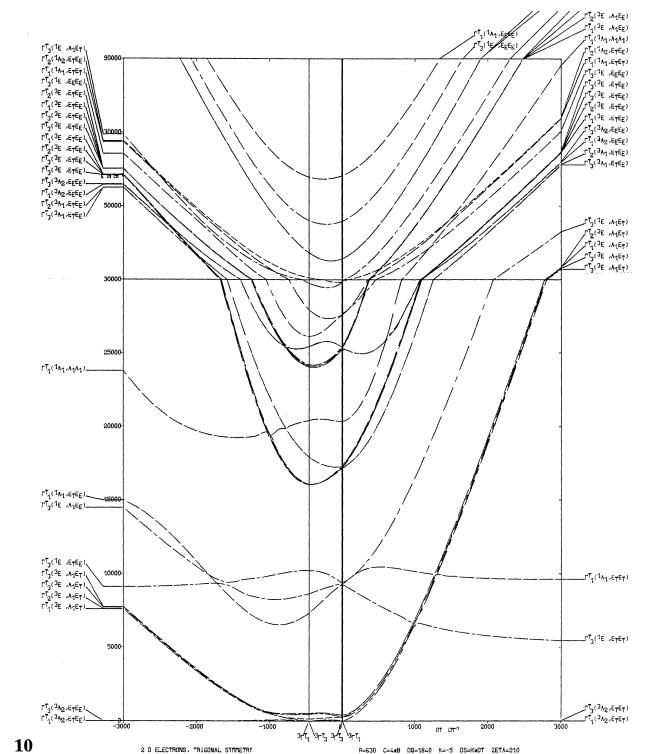
$O_h(T_d)$, no spin-orbit



(*)B Racah parameters related to the slater integrals (electronic repulsions)

- **König&Kremer (1970)**

Low symmetries+spin-orbit



$$A = F^0(3d,3d) - \frac{49}{441}F^4(3d,3d)$$

$$B = \frac{1}{49}F^2(3d,3d) - \frac{5}{441}F^4(3d,3d)$$

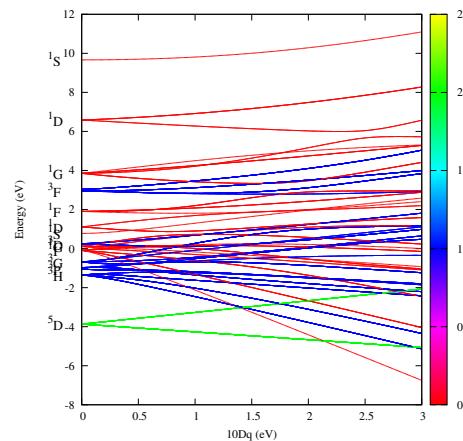
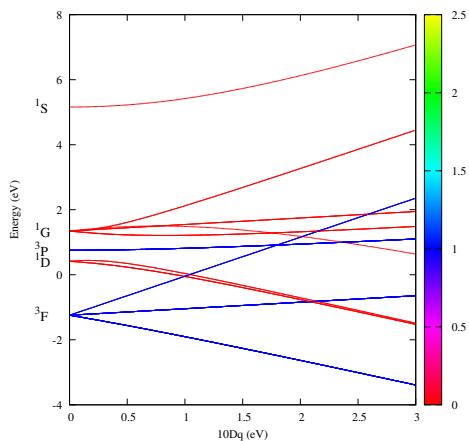
$$C = \frac{35}{441}F^4(3d,3d)$$

Crystal field : energy diagram

- Plot of the **energy** of spectroscopic terms ($^{2S+1}\Gamma$ or $^{2S+1}\Gamma_J$) as function of **crystal field parameter (10Dq,Ds,...)** (and B Racah parameter (*))
- **Tanabe-Sugano diagram (1954)**
 $O_h(T_d)$, no spin-orbit
- **König&Kremer (≈1970)**
Low symmetries+spin-orbit



Tutorial Monday afternoon :
02_Energy_Level_Diagram_Tanabe-Sugano.Quanty



Crystal field : multi-electron and orbitals

Multi-electron ions

- **Spectroscopic terms :** $\Gamma_i \in G$

Term written with **capital letter** in Mulliken notation. Ex : A_{1g}, T_2
or Koster notation : Γ_i

- **Electron density / orbitals**

ex : $e_g^{5.8} t_{2g}^{1.2}$

- **L,S,J,M_L,M_S,M_J not « good » quantum numbers.**

Expectation values:

$$\langle \Gamma_i | \hat{O} | \Gamma_i \rangle$$

with $\hat{O} = L_z, S_z, \dots$

One electron/orbitals

- **Orbitals :** $\gamma_i \in G$

Small letter

ex: e_g, t_{2g}

a_{1g}, e_2

Crystal field multiplet applied to core hole spectroscopies

Calculations including the core-hole (ex : $2p^53d^{n+1}$, $3d^94f^{n+1}$, ...)

Transitions : electric dipole, electric quadrupole, magnetic dipole...

- ✓ **Racah/Bander** (Theo Thole) : core of **CTM4XAS** (Frank de Groot)
Based on group theory (Butler “Point group symmetry applications”)
Difficult to modified the core in order to adapt to new spectroscopies since
Theo Thole died in 1996

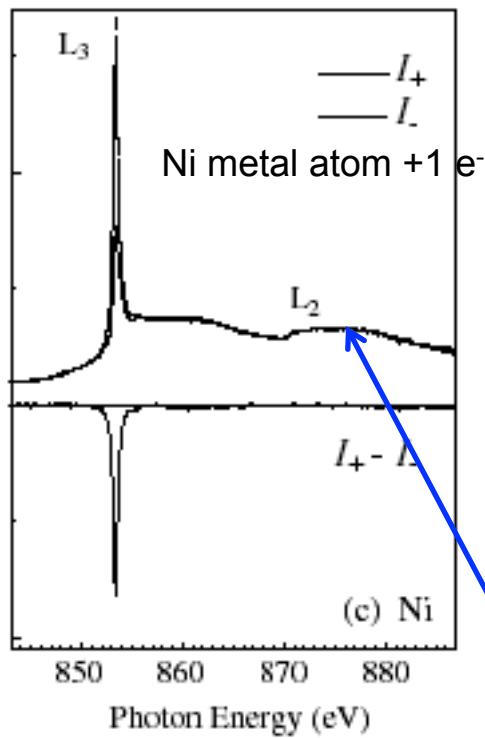
- ✓ **Quanty** (Mauris Haverkort) : core of **Crispy** (Marius Retegan)
and **CTM4RIXS** (Frank de Groot)
Flexible for the need of the new spectroscopies/physical properties

Example of crystal field effect

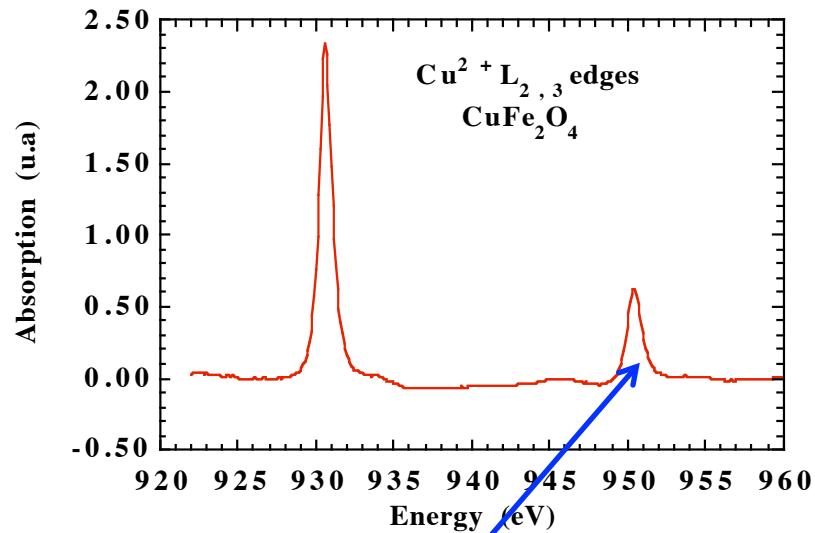
From spherical to Oh symmetry

*Example of crystal field effect:
 d^9 ion at $L_{2,3}$ edges ($2p \rightarrow 3d$)*

$3d^9$
 Spherical symmetry



$3d^9$
 $\approx O_h$ symmetry

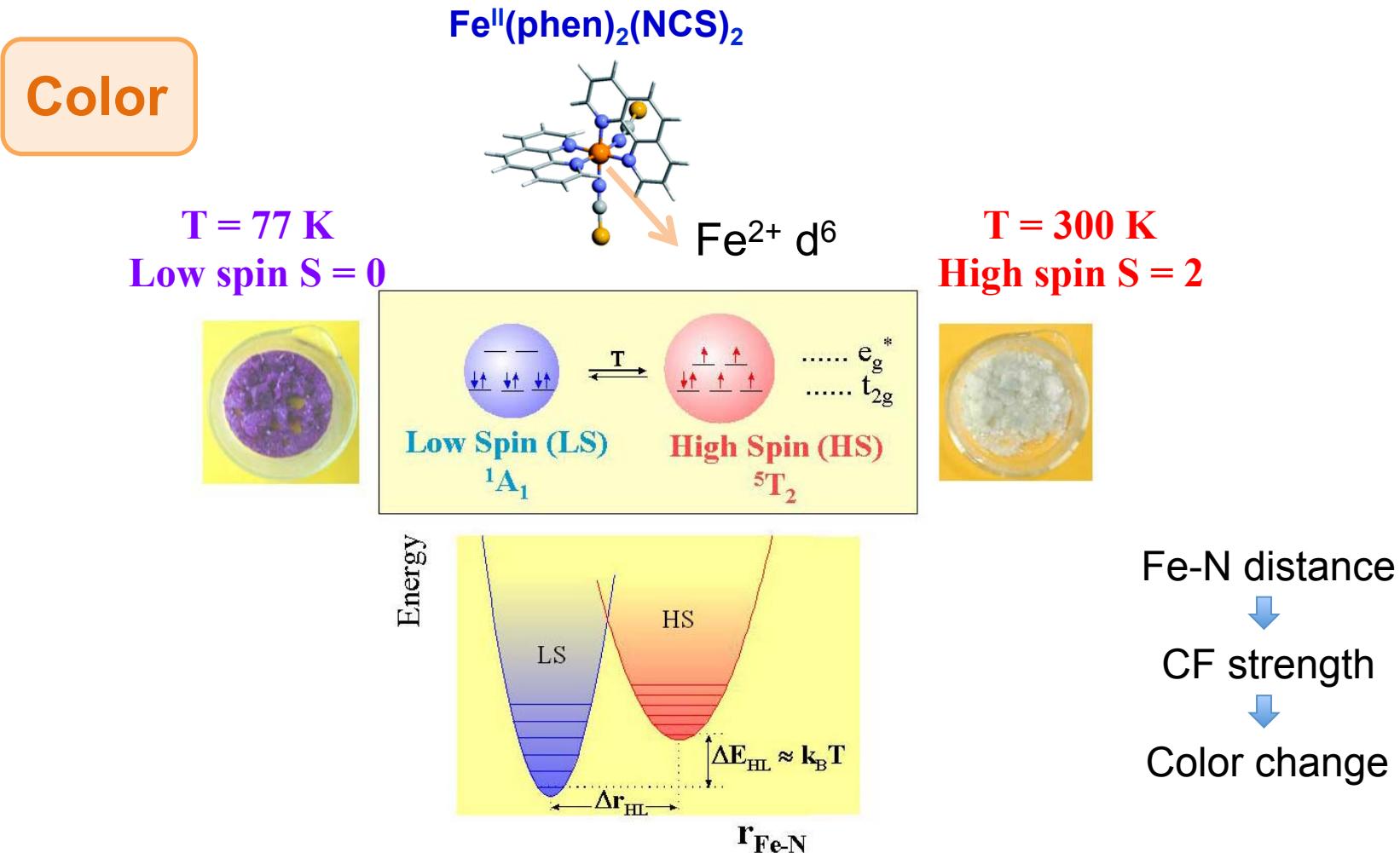


Can be understood via symmetry consideration using group theory

Example of crystal field effect

Spin crossover

Example of crystal field effect : spin crossover

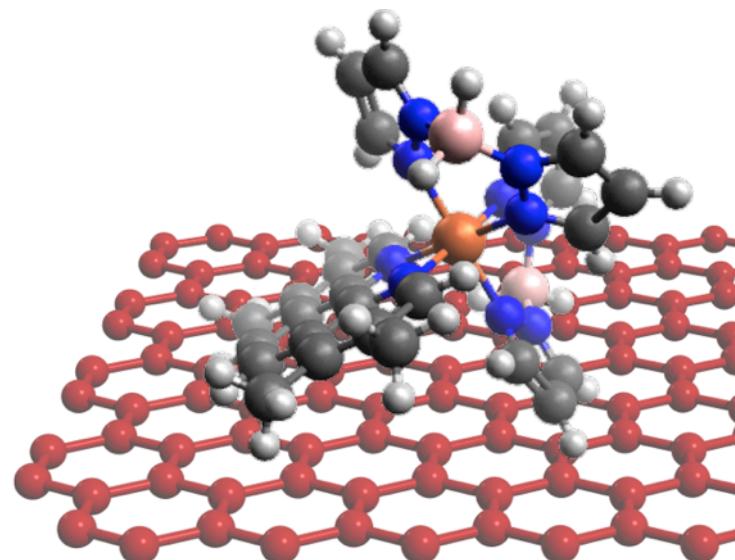
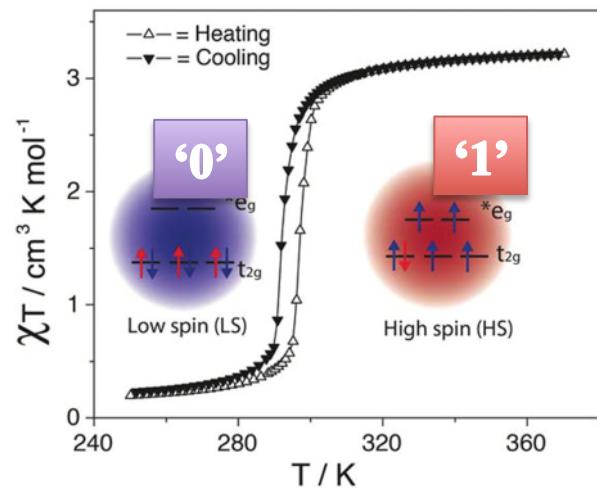


Ex of application: thermochromic painting

Example of crystal field effect : spin crossover

Magnetism:

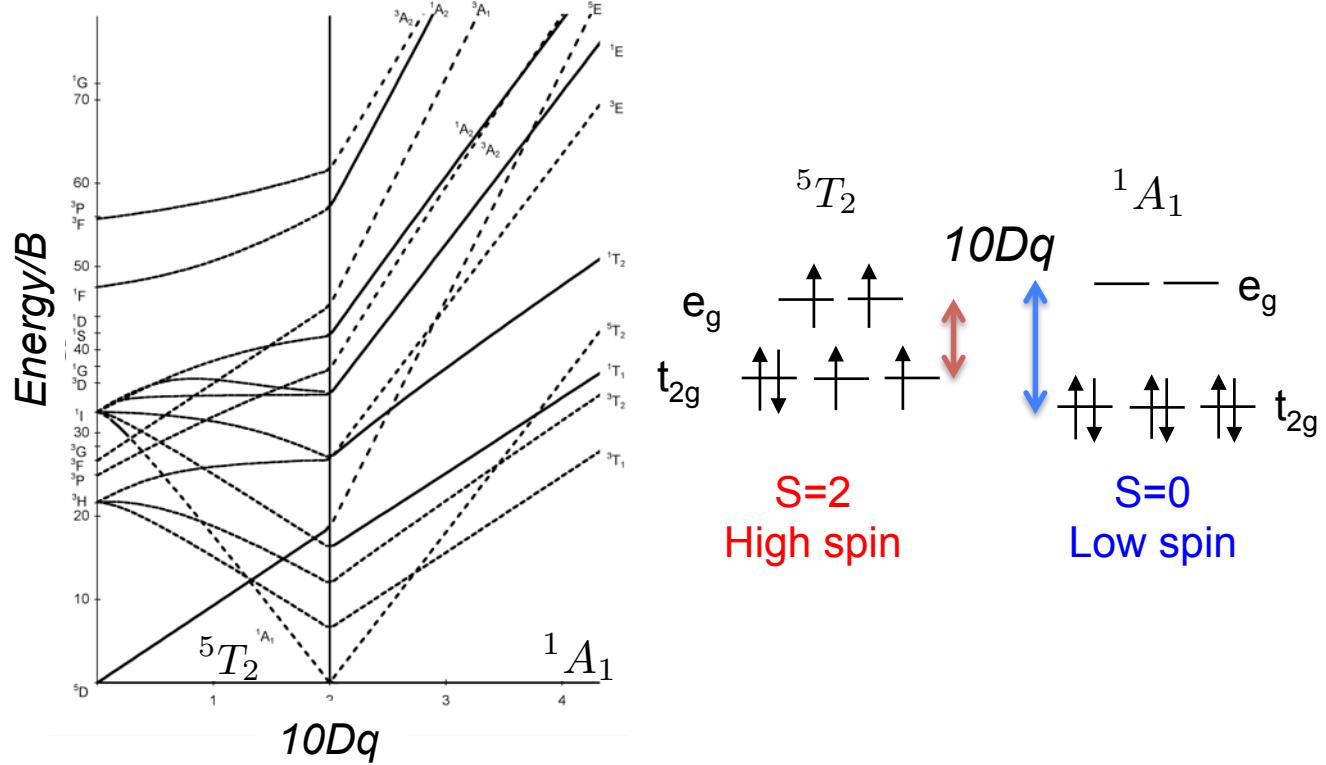
- switch between “0”(LS) and “1”(HS)
- driven by external stimuli (temperature, light, ...)
- couple to surface (insulating or conducting)
for **molecular spintronic**



Crystal field : spin crossover

Fe^{2+} ion ($3d^6$) in O_h symmetry

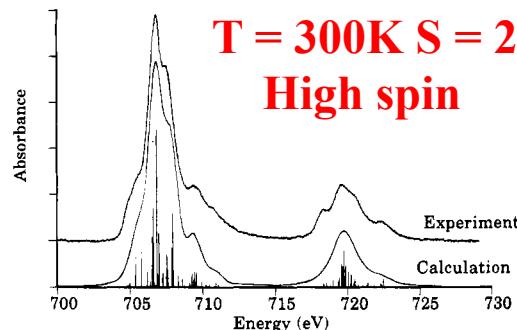
Temperature ↘
 Fe-N distance
 (observed by DRX)
 ↘ CF strength
 ($10Dq$)



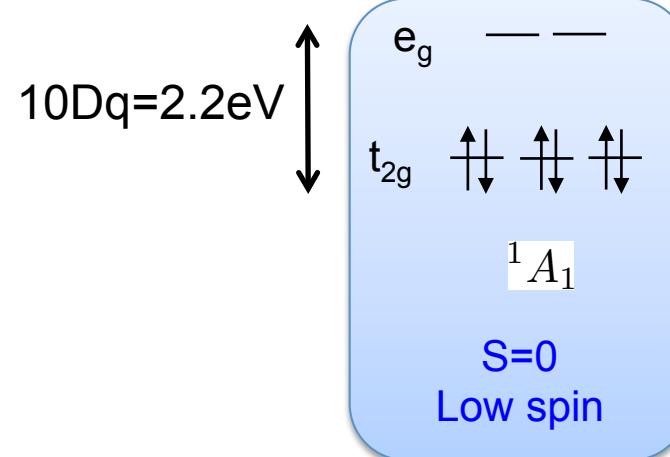
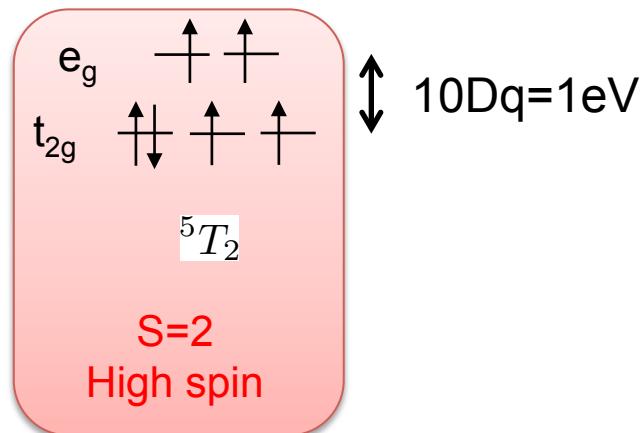
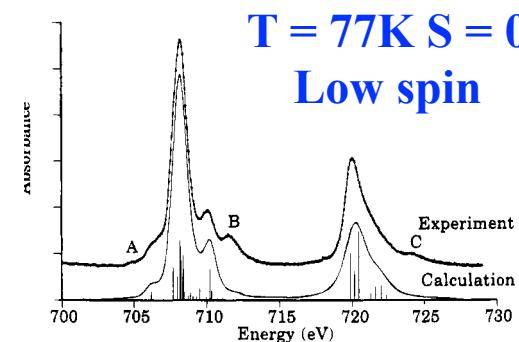
- Magnetic properties goes from paramagnetic to non-magnetic

Crystal field effect : spin crossover

XAS at Fe $L_{2,3}$ edges ($2p \rightarrow 3d$)



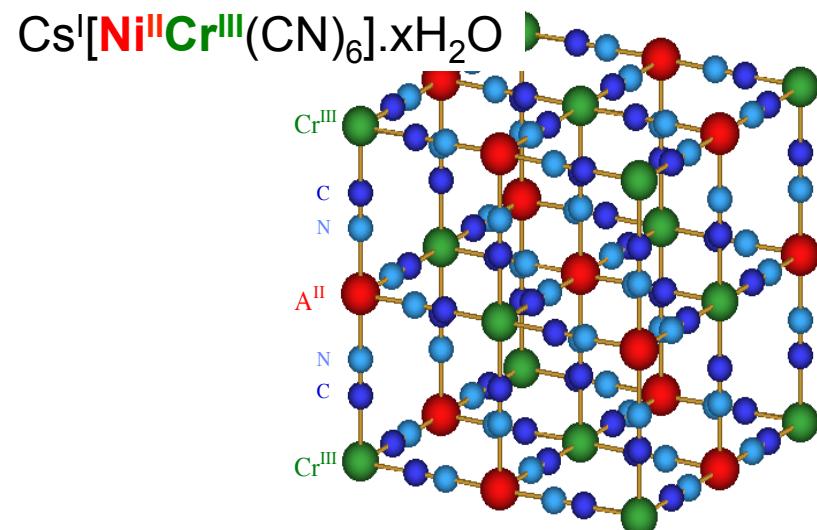
CFM calculation
 Fe^{2+} ($3d^6$)
 O_h symmetry



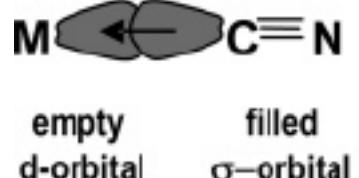
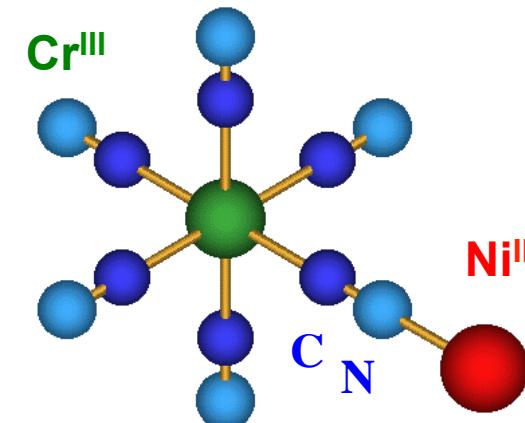
From Briois V., J. Am. Chem. Soc., 117 (1995)

Limit of the crystal field model: towards ligand field model

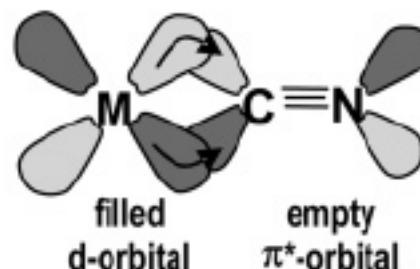
Example : prussian blue analogues



(covalent) $\text{Cr}^{\text{III}}\text{-C}\equiv\text{N-Ni}^{\text{II}}$ (ionic)



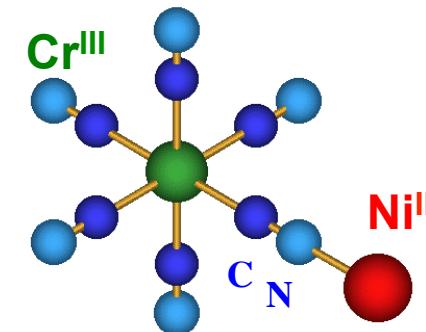
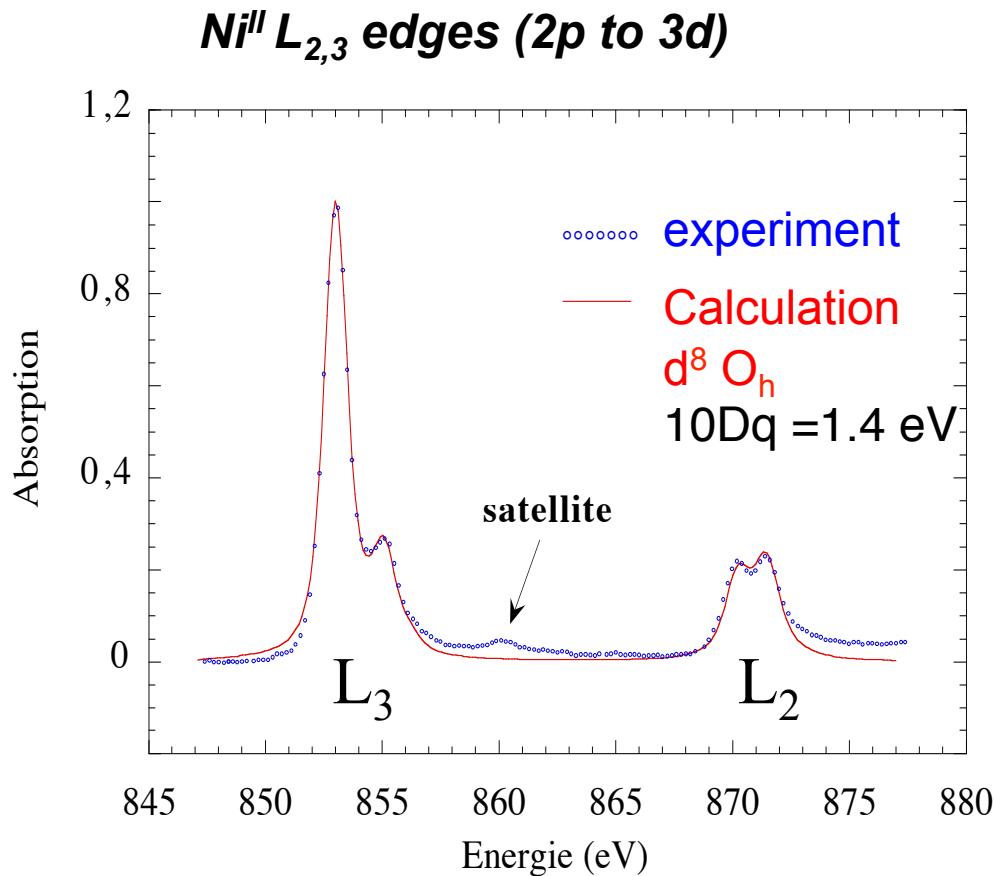
(i) σ donation



(ii) π back-donation

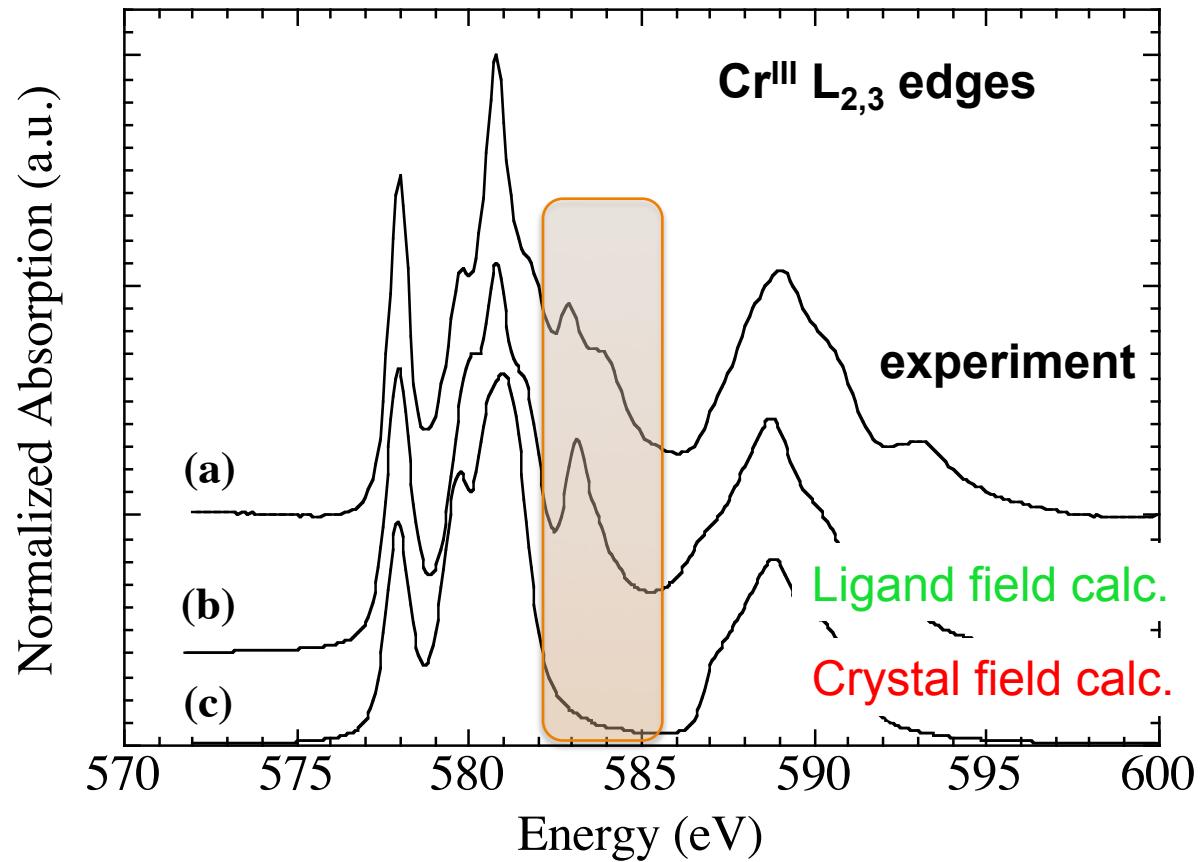
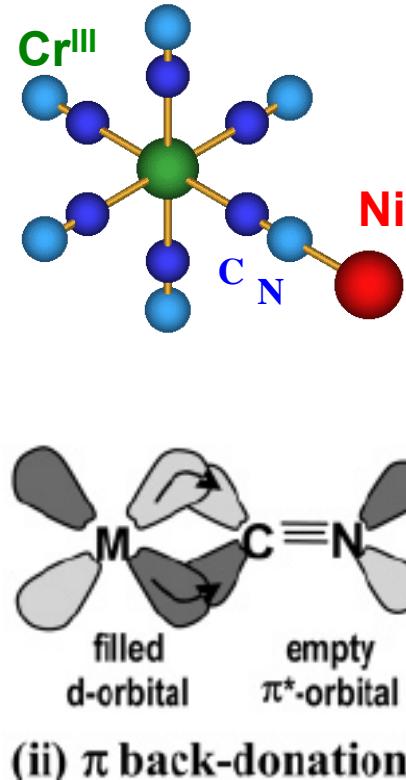
Weak covalent bond

CN-Ni^{II} in Cs[NiCr(CN)₆]



Weakness of crystal field theory

Strong crystal field

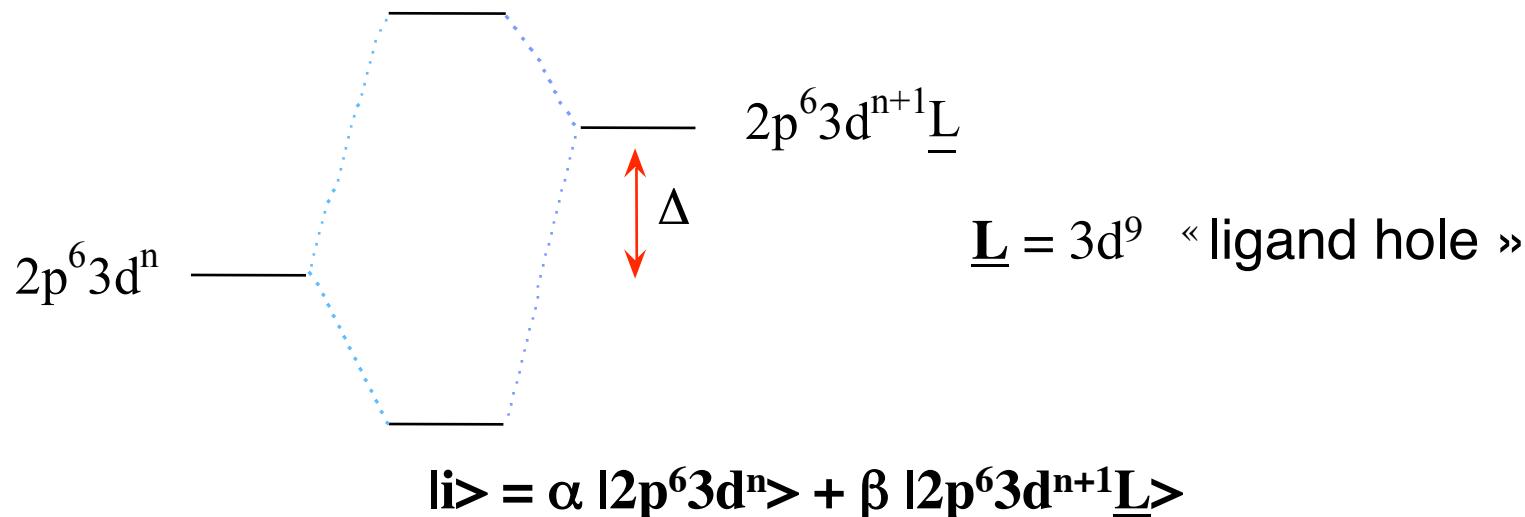


- **Crystal field** : Oh, large $10Dq$ (=3.5eV) + strong slater reduction (50%)
Incomplete model due to the strong metal to ligand charge transfer
- **Ligand field** : metal to ligand charge transfer
 $d^3d^0 + d^2d^1$

Ligand field multiplets

- To account for electron coming from ligands or neighbor metal
- **Multi-configurations** approach
 - also called cluster model
 - charge transfer model (CT)
- Two types of **charge transfer (CT)**:
 - Metal to ligand : **MLCT**
 - Ligand to Metal : **LMCT**

Ligand to metal charge transfer LMCT



Parameters

Δ energy (eV)

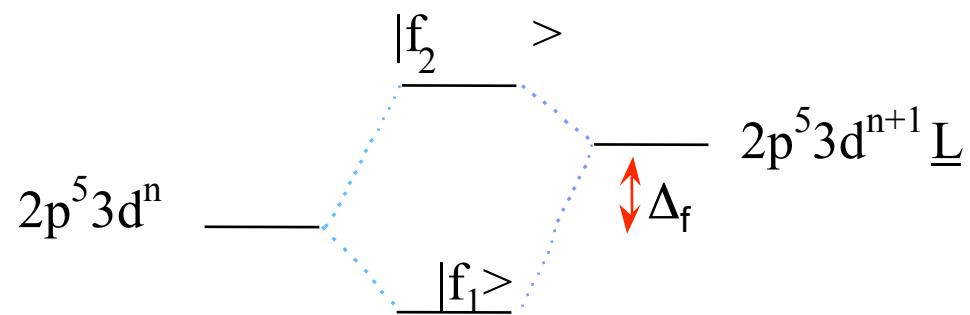
Charge transfer integral (symmetry dependent)

$$V = \langle 3d^n | H | 3d^{n+1}\underline{L} \rangle \text{ (also called } T\text{)}$$

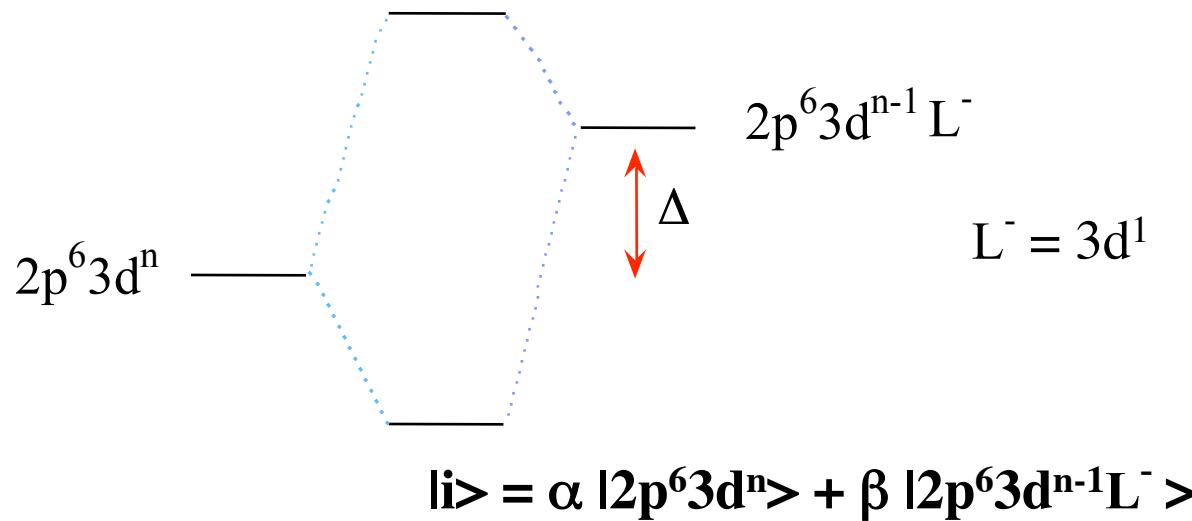
In O_h symmetry: $V_{t_{2g}}$ et V_{e_g}

Ligand to metal charge transfer LMCT

Core spectroscopies : excited state 2p for L_{2,3} edges
 1s for K edge



Metal to ligand charge transfer MLCT



Parameters

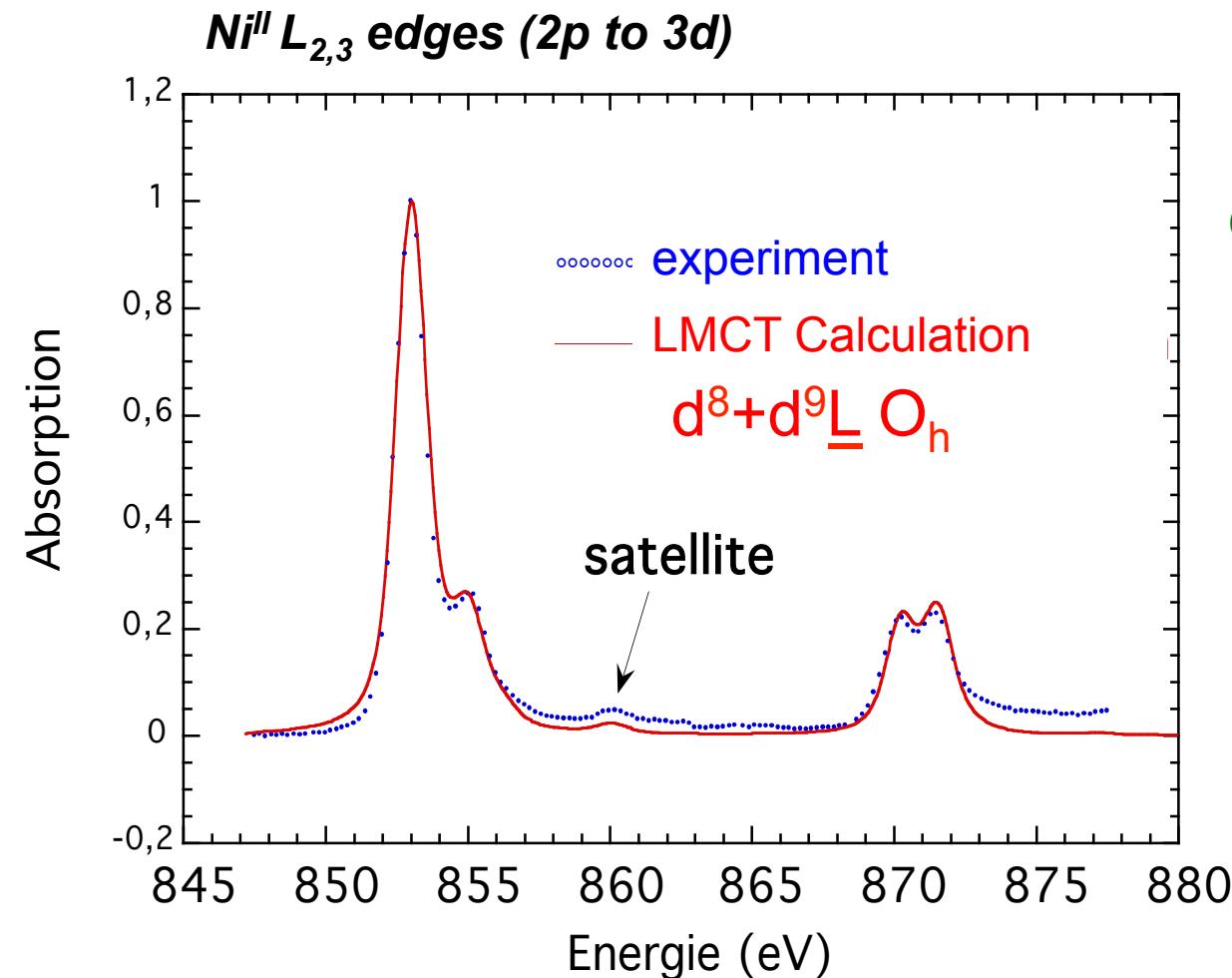
Δ energy (eV)

Charge transfer integral (symmetry dependent)

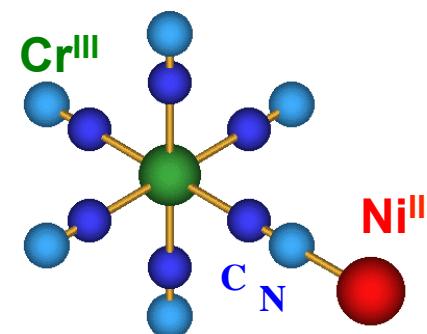
$$V = \langle 3d^n | H | 3d^{n-1}L^- \rangle$$

In O_h symmetry: $V_{t_{2g}}$ et V_{e_g}

LMCT

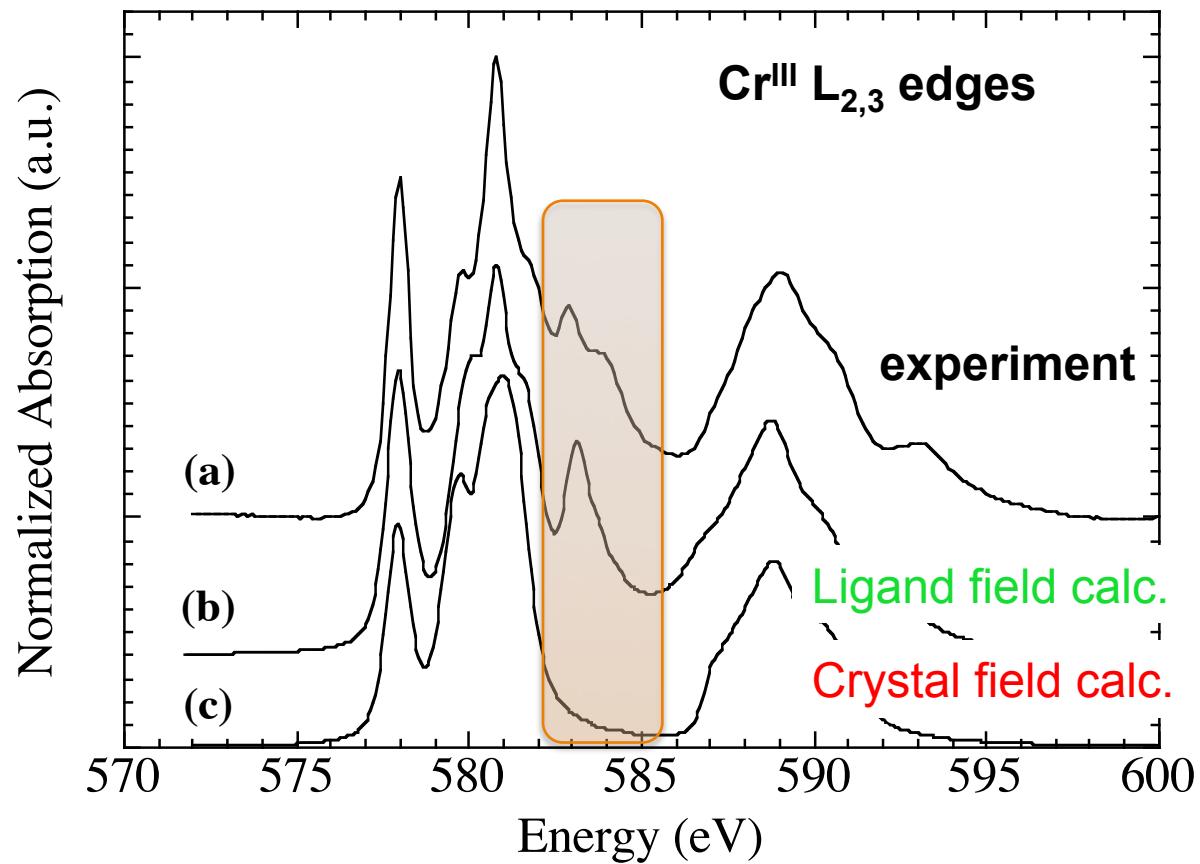
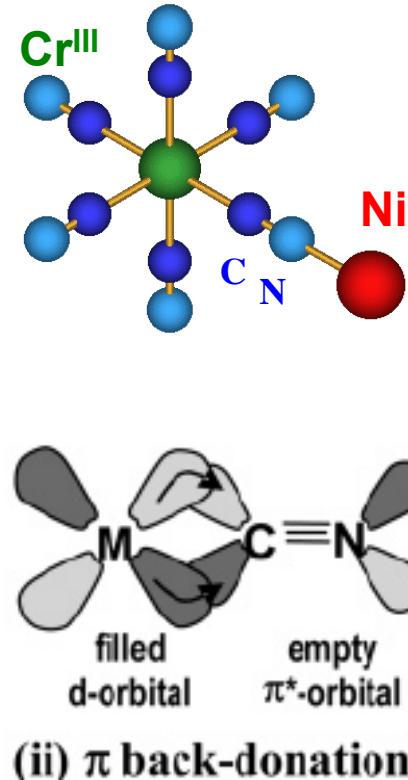


Cs[Ni^{II}Cr^{III}(CN)₆]



Weakness of crystal field theory

Strong crystal field



□ **Crystal field** : Oh, large $10Dq$ (=3.5eV) + strong slater reduction (50%)

□ **Ligand field** : MLCT $d^3+d^2L^-$

Crystal field and magnetism

1. Introduction / Definition
2. One electron (3d electron)
 1. Symmetries
 2. Hamiltonian
 3. Crystal field splitting of 3d electron
3. Multi-electrons ion (3d electrons)
 1. Hamiltonian
 2. Configurations/ spectroscopic terms
 3. Symmetry
 4. Examples

4. Crystal field and magnetism

1. Hamiltonian
2. Magnetic moment (spin and orbit)
3. Zeeman effect / paramagnetism

Magnetism

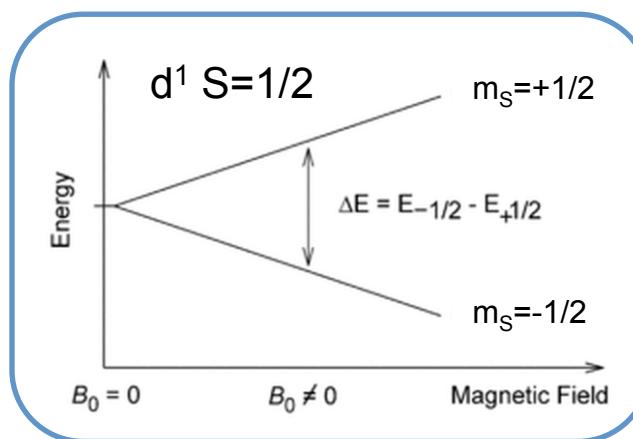
➤ Hamiltonian

$$H_{\text{ion}} = H_{\text{cin}} + H_{e-n} + H_{e-e} + H_{s-o} + H_{\text{CF}} \quad + H_{\text{Zeeman}} + H_{\text{Exchange}}$$

Free ion (spherical) Crystal field

$$\begin{aligned} H_{\text{Zeeman}} &= -\vec{m} \cdot \vec{H} = \mu_B (\hat{L} + g_0 \hat{S}) \cdot \vec{H} \\ &= \mu_B (L_x + 2S_x) H_x + \mu_B (L_y + 2S_y) H_y + \mu_B (L_z + 2S_z) H_z \end{aligned}$$

1 electron ($S=1/2$)
spherical



$$H_{\text{Exchange}} = S_x H_x + S_y H_y + S_z H_z$$

Magnetic moments

➤ Definitions

$$m = m_{\text{orbit}} + m_{\text{spin}}$$

$$m_{\text{orbit}} = -\langle i | \hat{L}_z | i \rangle \mu_B$$

$$m_{\text{spin}} = -g_o \langle i | \hat{S}_z | i \rangle \mu_B = -2 \langle i | \hat{S}_z | i \rangle \mu_B$$

➤ Crystal field effect

$|i\rangle = |(L, S)J\Gamma\gamma\rangle$ are not eigenfunctions of \hat{L}_z or \hat{S}_z

$$\langle i | \hat{L}_z | i \rangle \neq m_L$$

$$\langle i | \hat{S}_z | i \rangle \neq m_S$$

$H_{\text{cin}} + H_{\text{e-n}} + H_{\text{e-e}} + H_{\text{s-o}}$ commute with $\hat{J}^2, \hat{J}_z, \hat{L}^2, \hat{S}^2$
 $(H_{\text{cin}} + H_{\text{e-n}} + H_{\text{e-e}} \text{ commute with } \hat{L}^2, \hat{L}_z, \hat{S}^2, \hat{S}_z)$

Magnetic moments and spectroscopies

- For 3d ions, the orbital magnetic moment is quenched (nul) in most cases.

Some 3d ions have significant orbital magnetic moment
(Co²⁺, low spin Fe³⁺, Fe²⁺)

→ **magnetic anisotropy**

- **Magnetic core spectroscopies**

XMCD, RIXS-MCD, XMLD, XM χ D

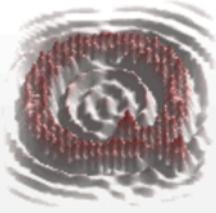
→ **access to**

Spin and orbit magnetic moment
Anisotropies + Angular dependences

Magnetic moments in Quanty

➤ **Quanty**: any operator (\hat{O}) and its expectation values

$$\langle i | \hat{O} | i \rangle$$

 Ex : spin operator

Article [Talk](#) Read [Show pagesource](#)

S_z

The S_z operator is defined as:

$$S_z = \sum_{\tau} \frac{1}{2} (a_{\tau\uparrow}^\dagger a_{\tau\uparrow} - a_{\tau\downarrow}^\dagger a_{\tau\downarrow}).$$

The equivalent operator in Quanty is created by:

Example.Quanty

```
OppSz = NewOperator("Sz", NF, IndexUp, IndexDn)
```

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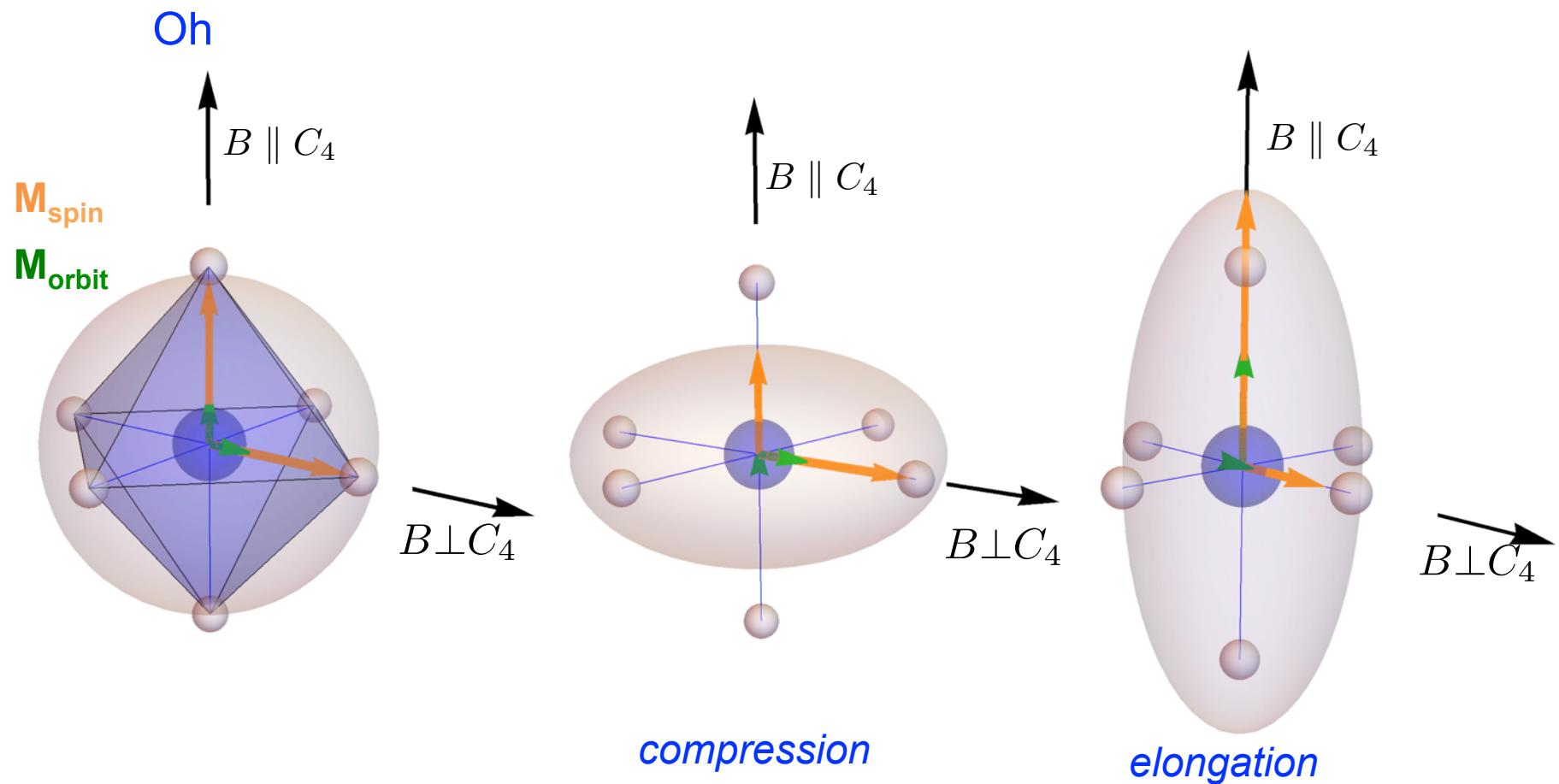
- Smin
- Splus
- Ssqr
- Sx
- Sy
- Sz



**Tutorial 07_Expectationvalues.Quanty
(Monday morning)**

Crystal field and magnetism

Tetragonal distortion : D_{4h}



Conclusion / remarks

- Crystal field and ligand field are atomic models
- Crystal/ligand field are a semi-empirical model (parameter dependent)
can be linked to more *ab-initio* methods (tight-binding, DFT, CASSCF,...)
- Crystal field can create **natural and/or magnetic anisotropies**
angular dependences
measured by core spectroscopies using polarized X-rays (linear or circular)



Following Lectures

- Crystal field affect the properties of **4f rare earth ions** (not discussed here)
More parameters